

Stability of PAH clusters

Sébastien Zamith, Ming-Chao Ji, Jean-Marc L'Hermite, Christine Joblin,

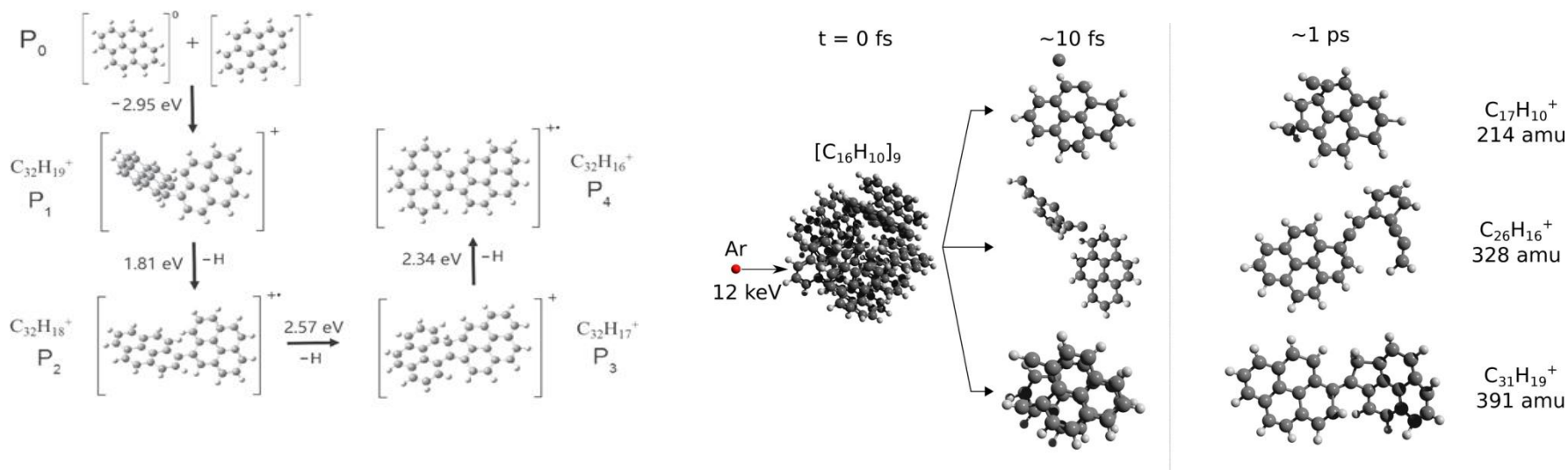
Léo Dontot, Mathias Rapacioli and Fernand Spiegelman



Stability of PAH clusters

PAH clusters as models for carbonaceous nanograins.

Destruction of such nanograins: route for formation of large complex PAHs?



Zhen et al. *ApJ* 863 128 (2018)

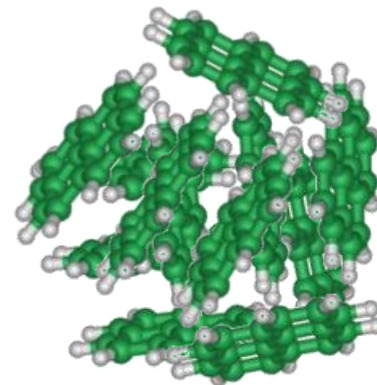
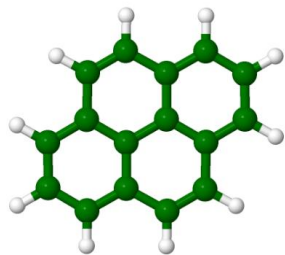
Delaunay et al, *J. Phys. Chem. Lett.*, 6, 1536-1542 (2015)

Are these PAH clusters stable?

Stability of PAH clusters

- Cationic pyrene clusters $(\text{Py})_n^+$
studied using
- Mass spectrometry techniques
and
- Phase Space Theory
to deduce
- Dissociation energies

Pyrene : $\text{C}_{16}\text{H}_{10}$

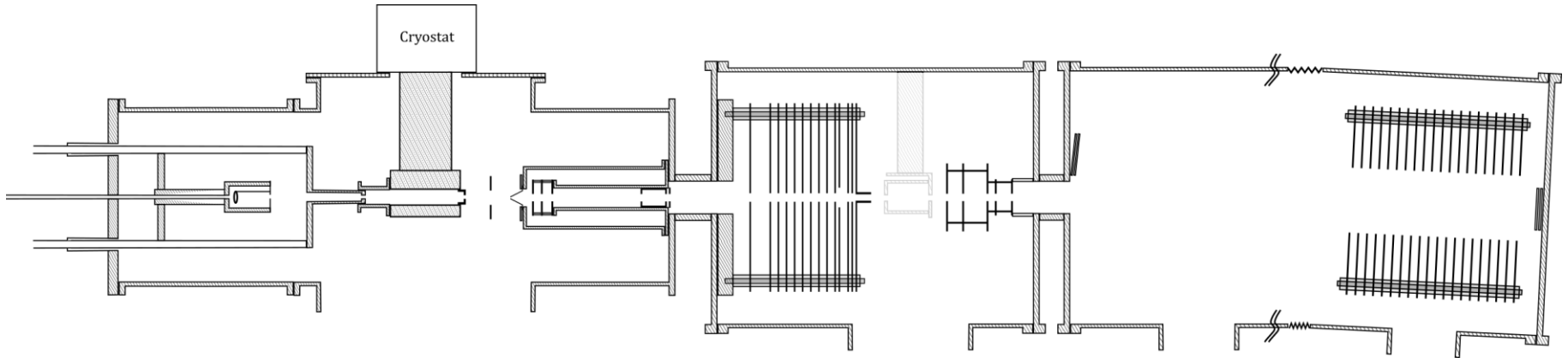


PAH = polyaromatic hydrocarbons

Outline

- Experimental setup
- Experimental results
- Phase Space Theory
- Dissociation energies
- Conclusion

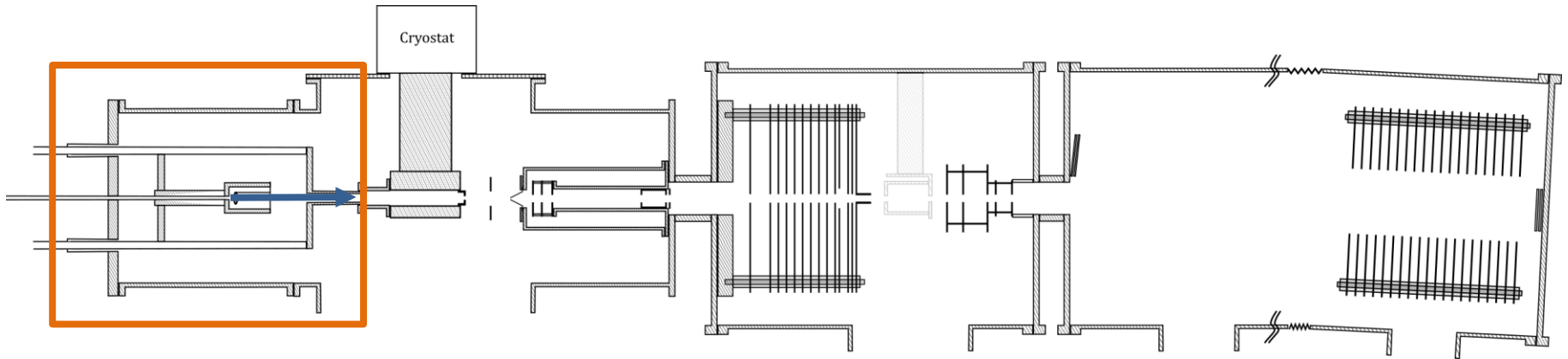
Experimental setup



Experimental setup initially designed to perform collisions between mass selected clusters and atomic or molecular vapor (attachment cross-section, fragmentation cross-section, nanocalorimetry, ...).

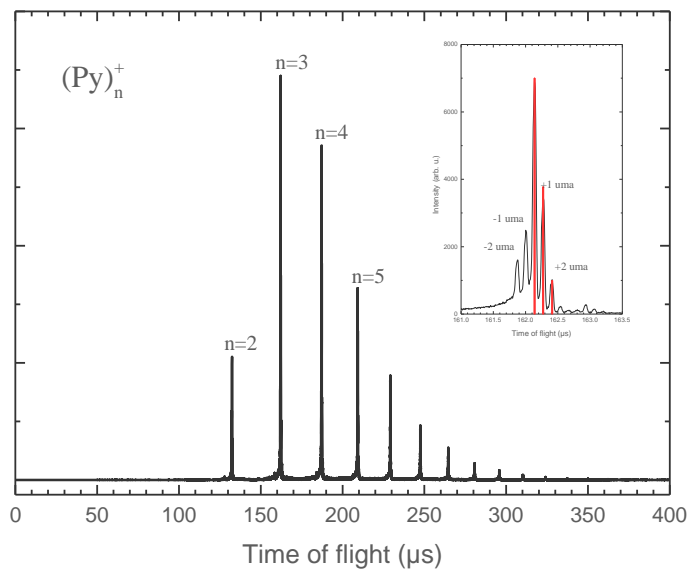
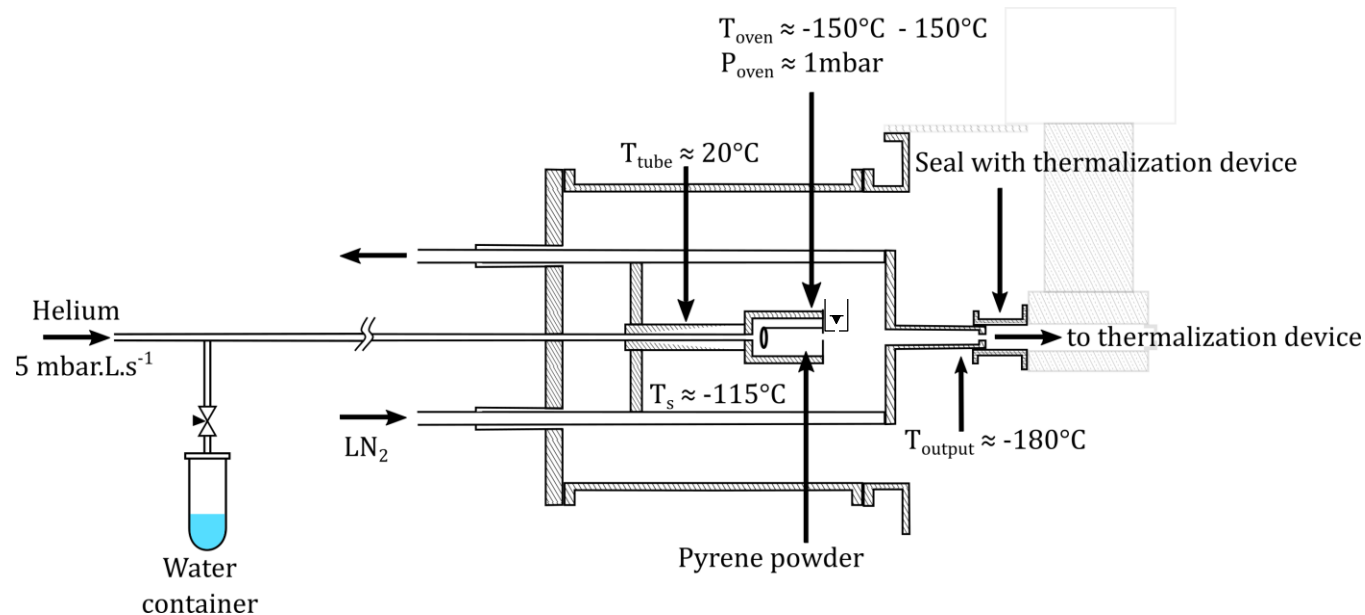
Can also be used to observe the **spontaneous thermal evaporation of mass selected clusters**.

Experimental setup



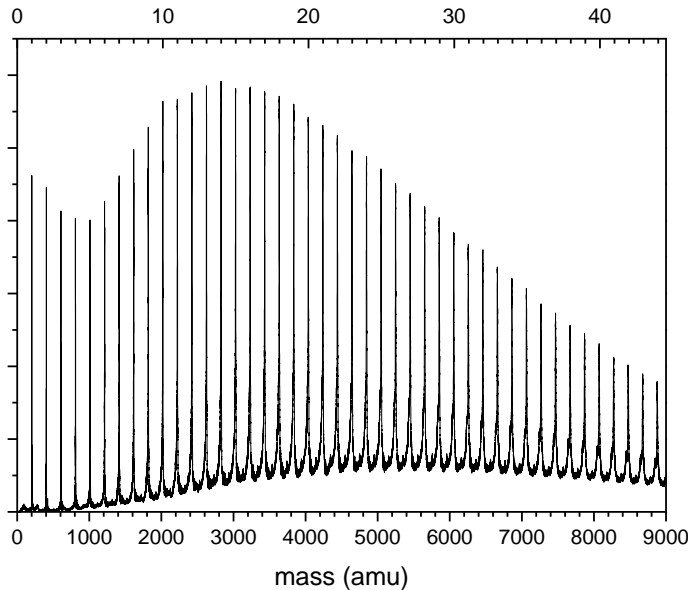
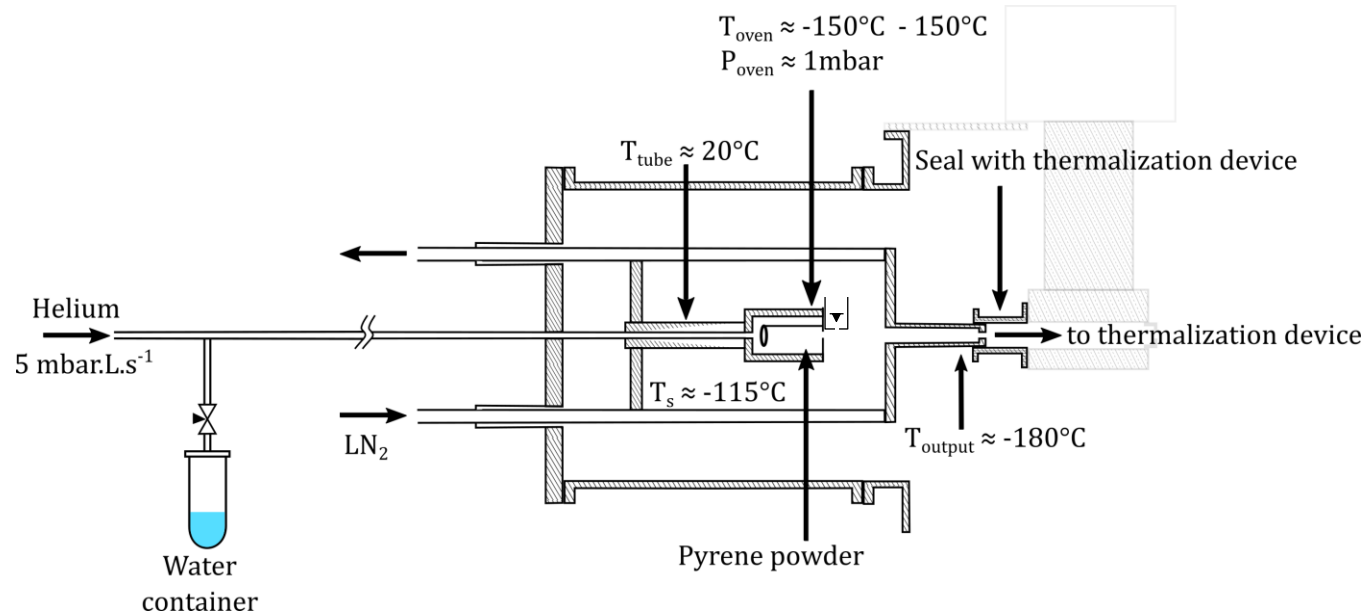
- Gaz aggregation source

Cluster production



Small pure Pyrene clusters

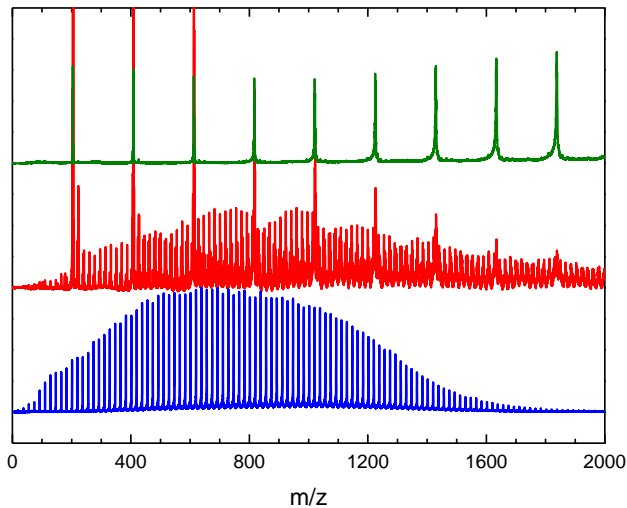
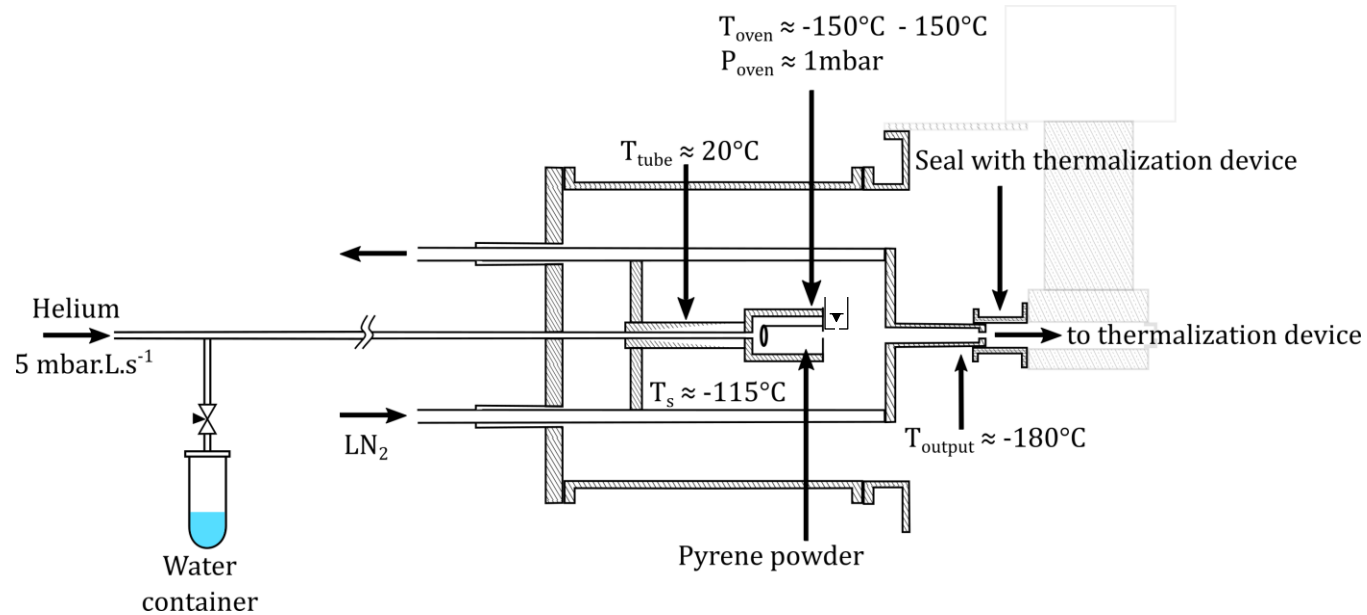
Cluster production



Small pure Pyrene clusters

Large pure pyrene clusters

Cluster production



Small pure Pyrene clusters

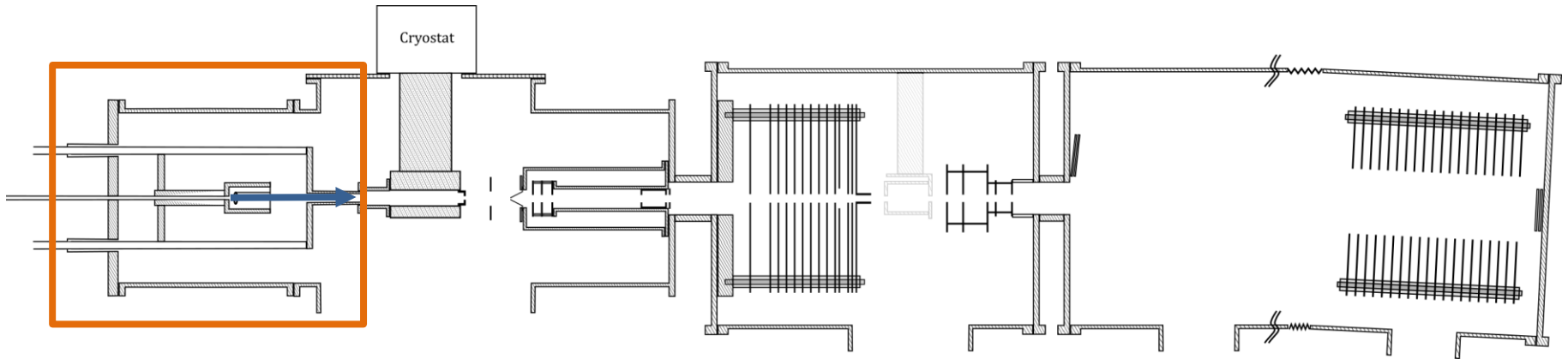


Large pure pyrene clusters

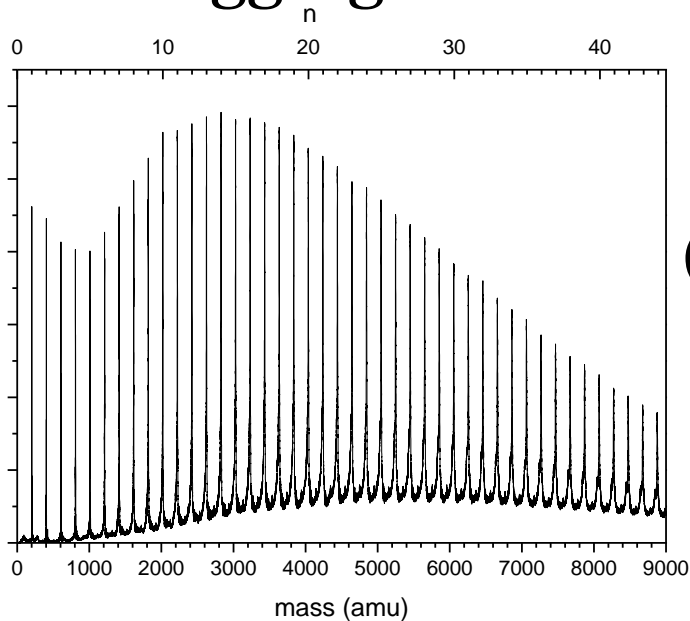


Mixed water-pyrene clusters

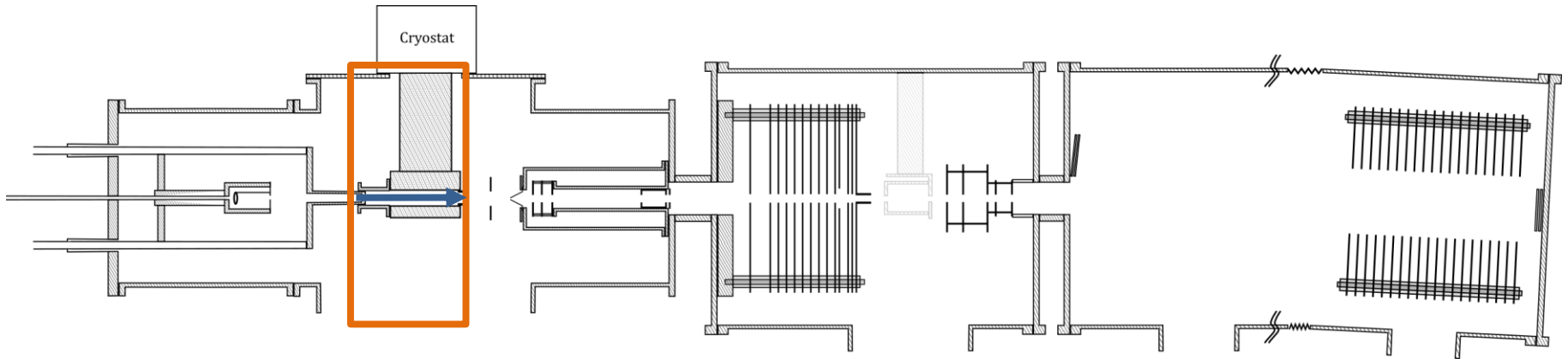
Experimental setup



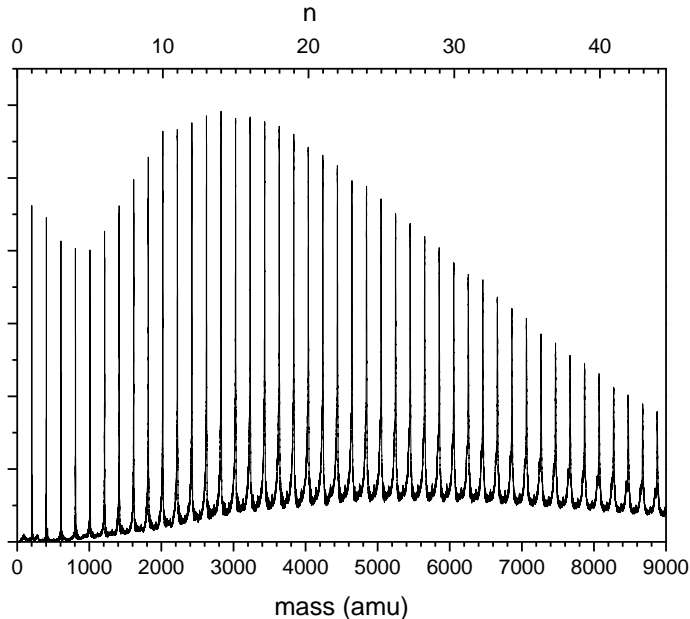
- Gaz aggregation source, $n = 1-40$



Experimental setup



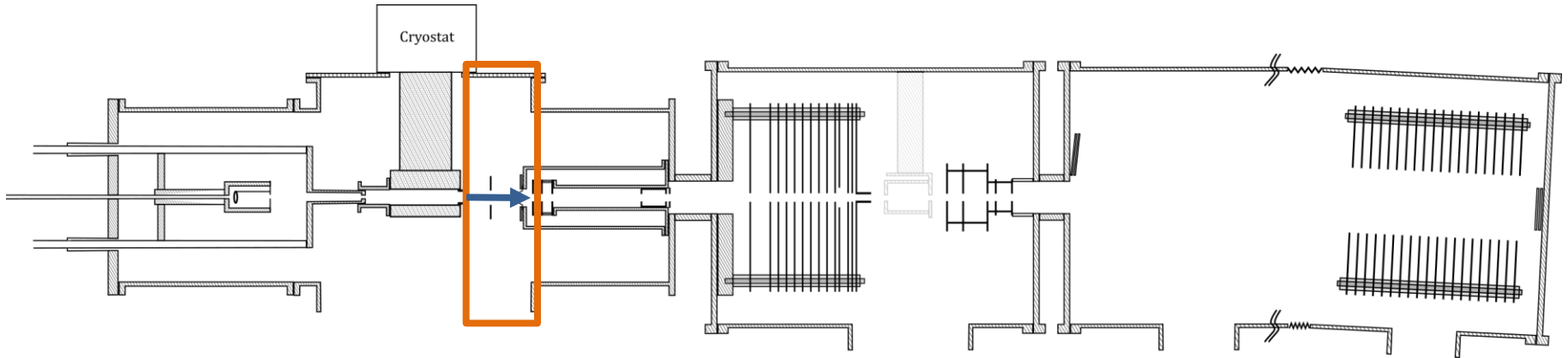
- Thermalization : $T = 25-300$ K



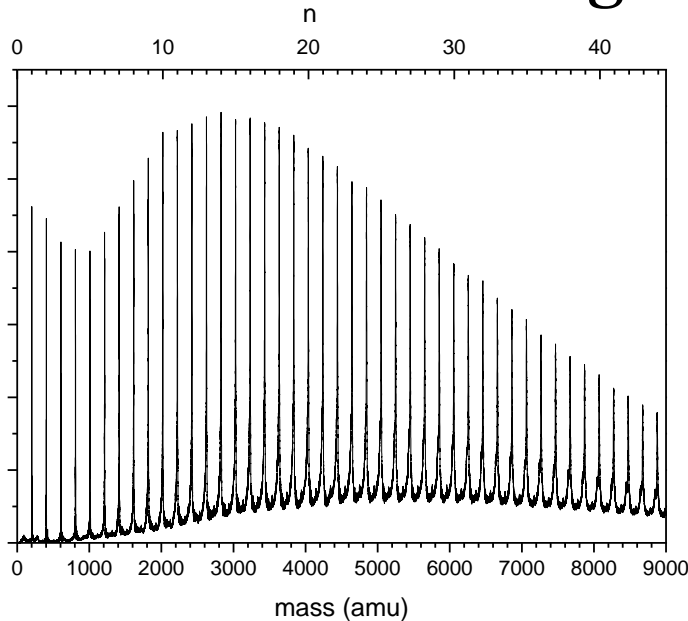
Large number of collisions
with the helium buffer gas

Canonical distribution
of internal energies E_i

Experimental setup

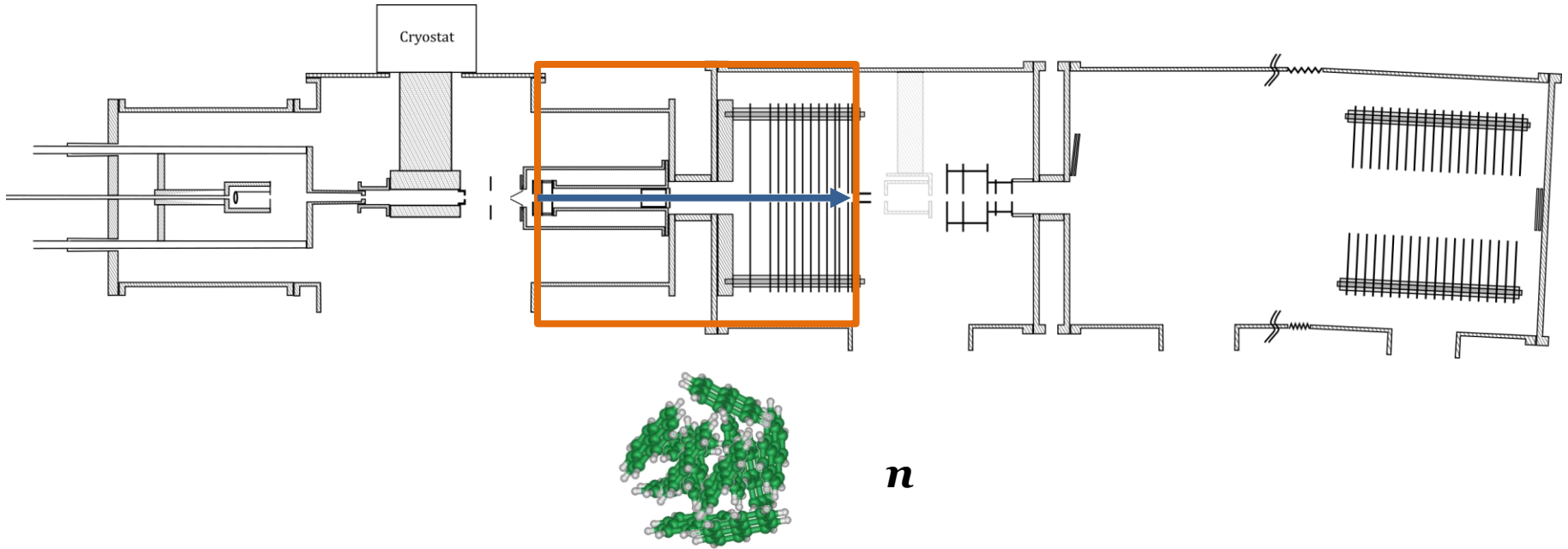


- Transfer to the high vacuum part



Internal energies E_i , microcanonical evolution

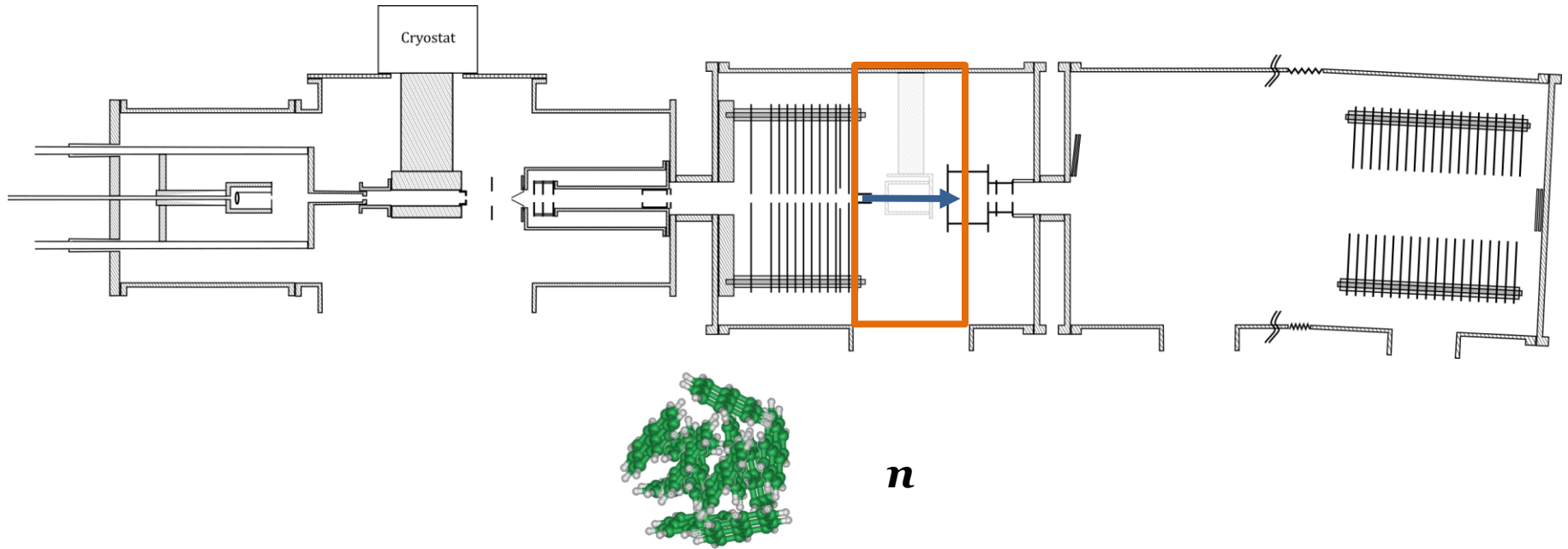
Experimental setup



- Mass selection and slowing down

Internal energies E_i , microcanonical evolution

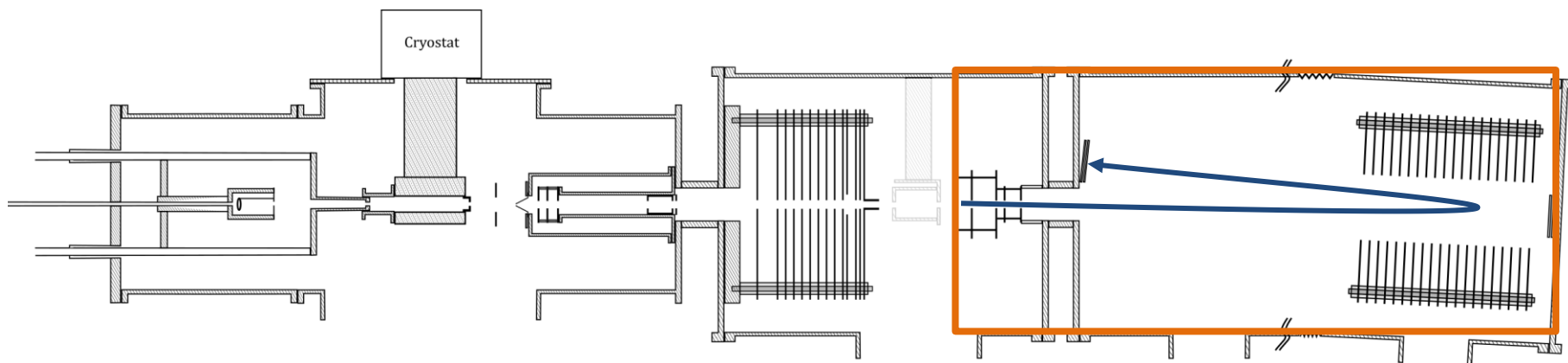
Experimental setup



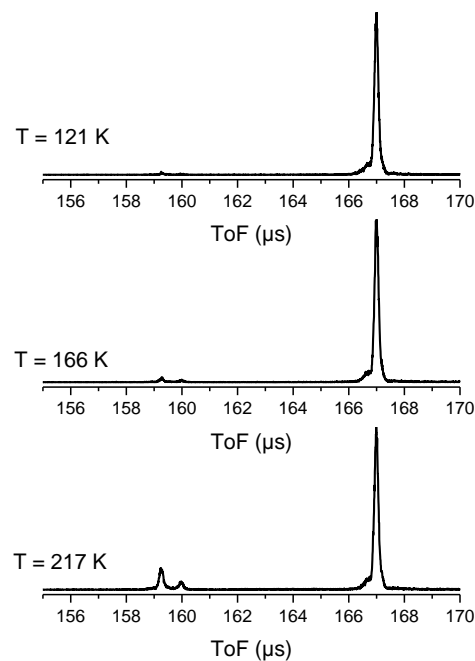
- Free flight

Internal energies E_i , microcanonical evolution

Experimental setup



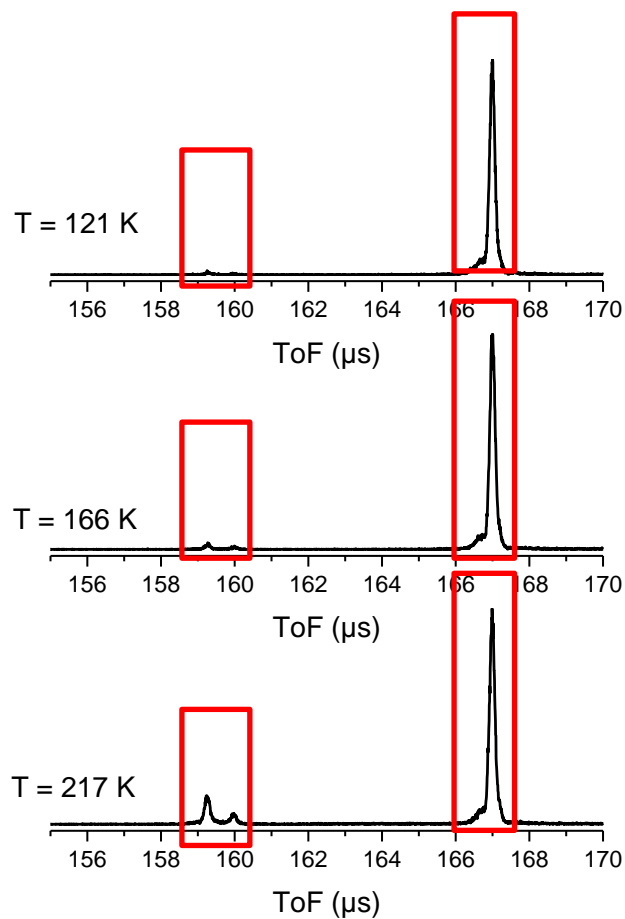
- Products analyzed by TOF mass spectrometry



Experimental results

Evolution of the mass spectra of mass selected clusters with initial temperature.

As the temperature is raised, appearance of fragments due to evaporation



Example: $(\text{Py})_{11}^+$ @ 22 eV

Ratio I/I_0

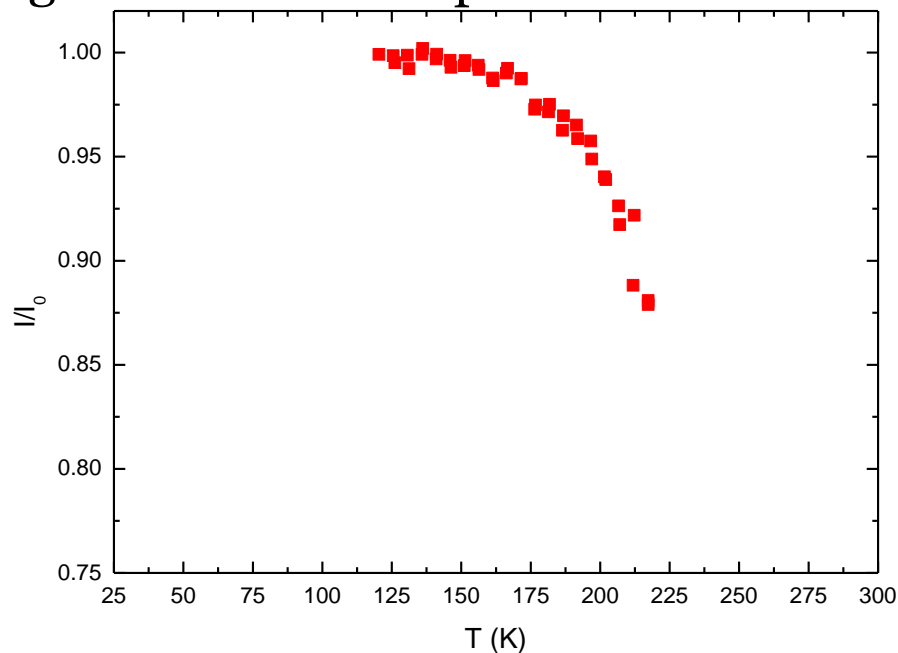
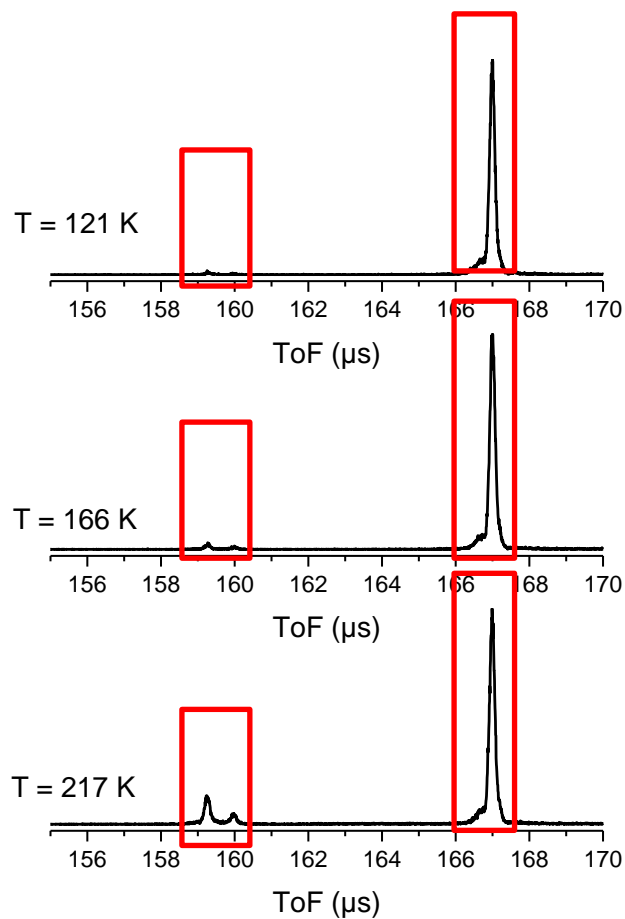
I = parent peak intensity

I_0 = sum of parent + fragment peaks

Experimental results

Evolution of the mass spectra of mass selected clusters with initial temperature.

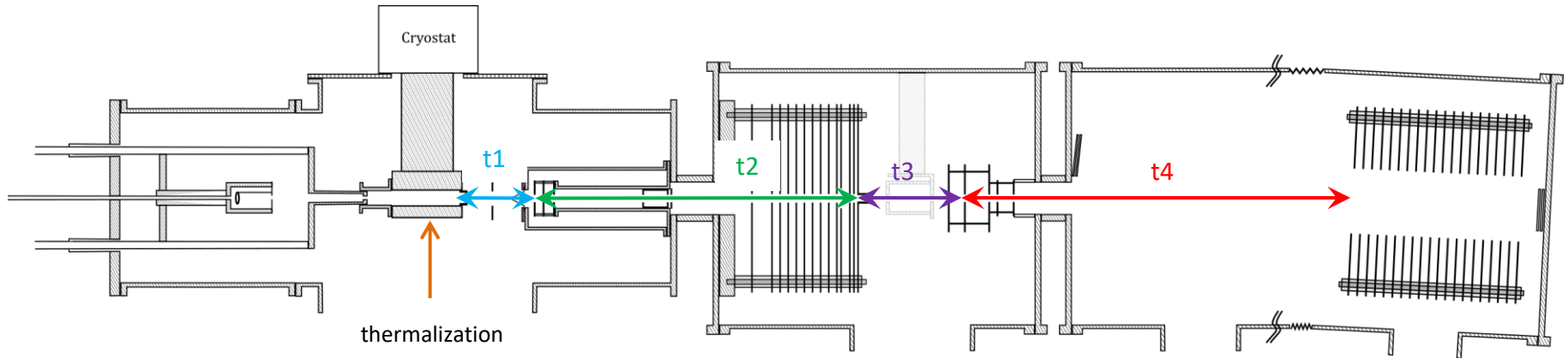
As the temperature is raised, appearance of fragments due to evaporation



Example: $(\text{Py})_{11}^+$ @ 22 eV

$$I = I_0 e^{-W(T)t} ?$$

Experimental results



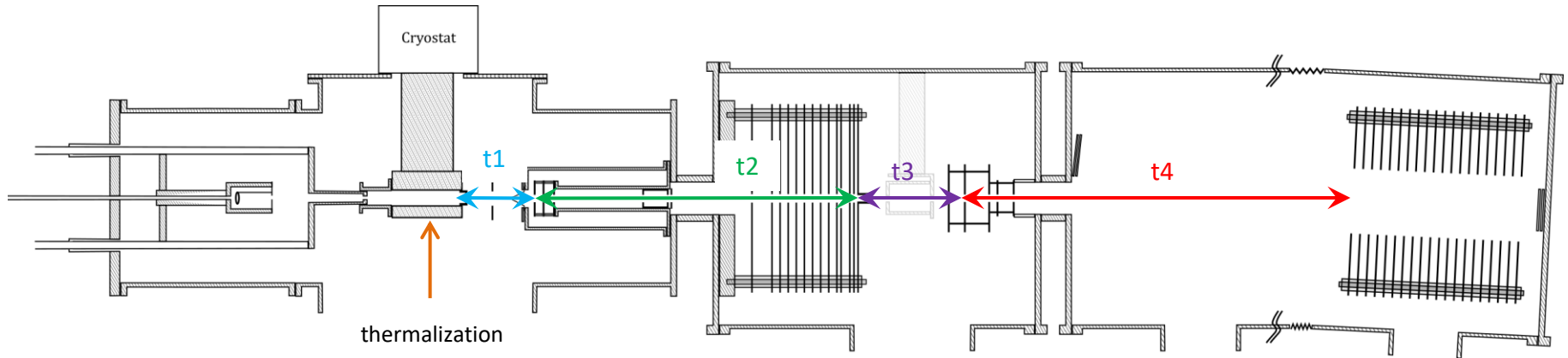
At the thermalizer exit:

- Cluster size distribution
- Canonical distribution of internal energies E_i

After the propagation time t_1 , the population of size n will have contributions from the evaporation of larger sizes:

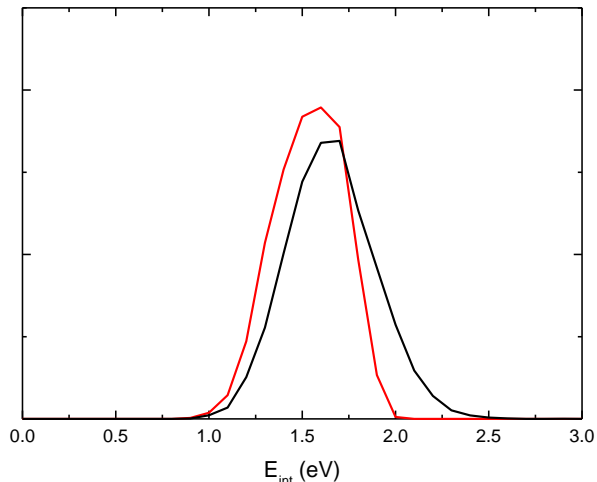
$$n + 1, E_i \rightarrow n, E'_i$$

Experimental results



At the thermalizer exit:

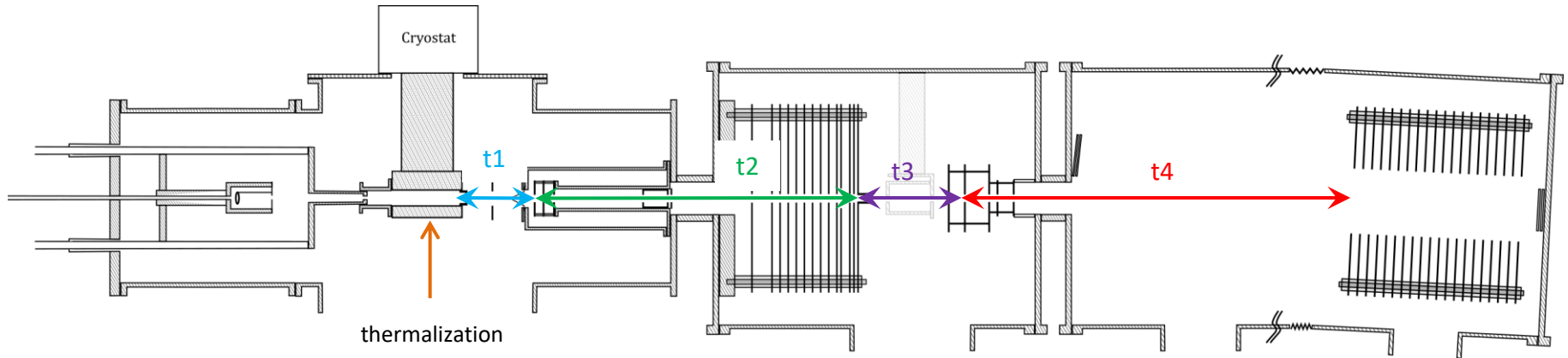
- Cluster size distribution
- Canonical distribution of internal energies E_i



After t_1 , population of size n might no longer be at temperature T

⇒ **new**, non-canonical, distribution of internal energies

Experimental results



At the thermalizer exit:

- Cluster size distribution
- Canonical distribution of internal energies E_i

Depending on where evaporation takes place, it might not be observed.

Experimental results reproduced by simulating the propagation in the setup with evaporation probabilities evaluated at each time step

⇒ Model for evaporation rates

PST evaporation rates

Ingredients of the Phase Space Theory:

- initial internal energy of the parent E_i
- density of states of the parent $N(E)$
- total number of states of the fragments $G(E, J)$
- conservation of angular momentum

$$W(E_i, J) = \frac{G(E_f, J)}{h(2J + 1)N(E_i)}$$

$$E_f = E_i + E_{rot} - D$$

$$E_{rot} = B_0 J(J + 1)$$

Approximations:

- only vibrational harmonic frequencies considered
- all species considered as spherical tops
- ion-polar interaction between the neutral fragment and the charged cluster

Energy partition among the fragments

Harmonic frequencies and moments of inertia from DFTB calculation (cf Rapacioli *et al.*)

PST evaporation rates

Ingredients of the Phase Space Theory:

- initial internal energy of the parent E_i
- density of states of the parent $N(E)$
- total number of states of the fragments $G(E, J)$
- conservation of angular momentum

Only one adjustable parameter

$$W(E_i, J) = \frac{G(E_f, J)}{h(2J + 1)N(E_i)}$$

$$E_f = E_i + E_{rot} - D$$
$$E_{rot} = B_0 J(J + 1)$$

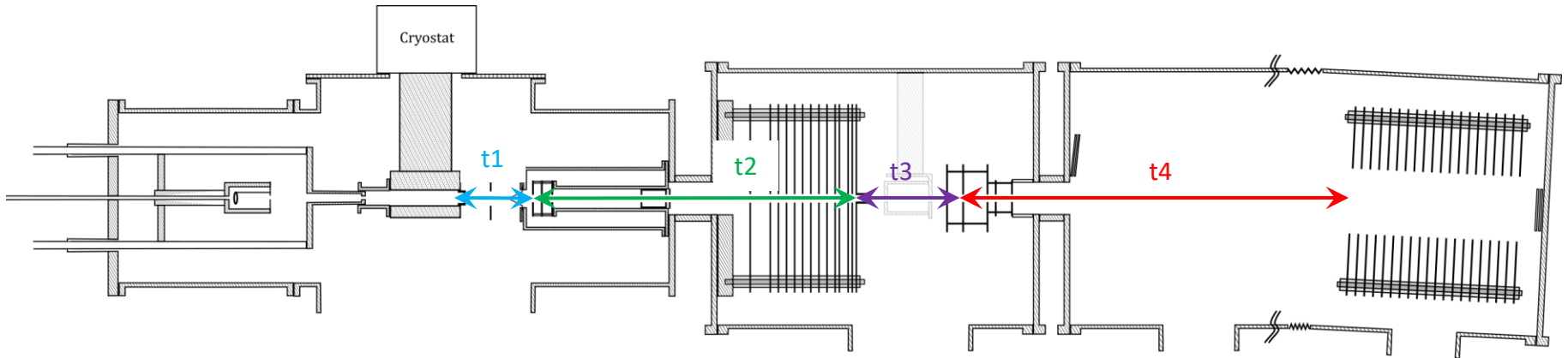
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PST evaporation rates

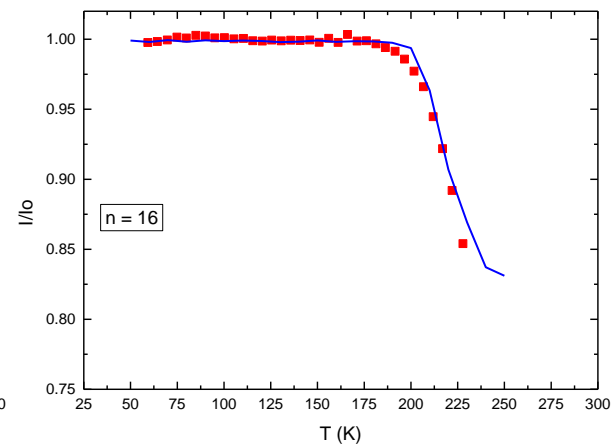
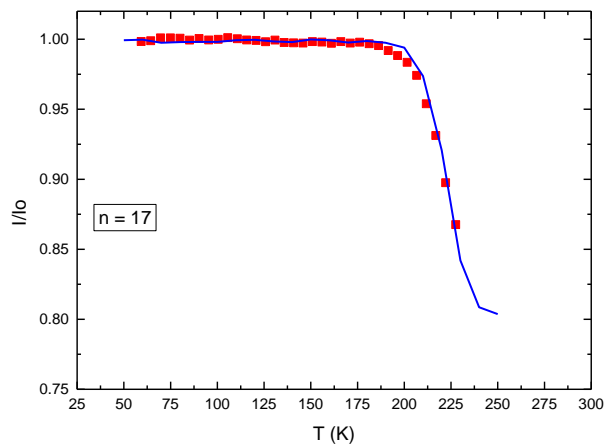
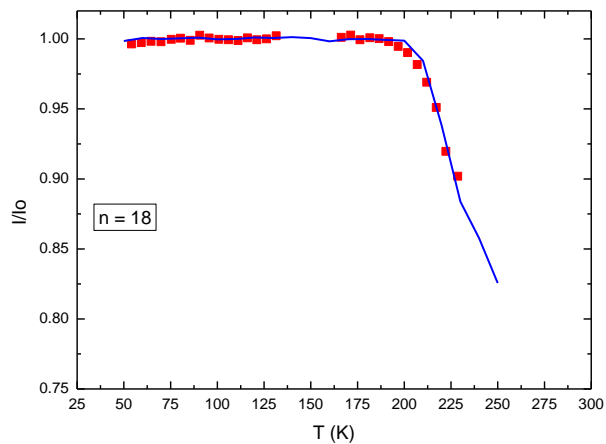
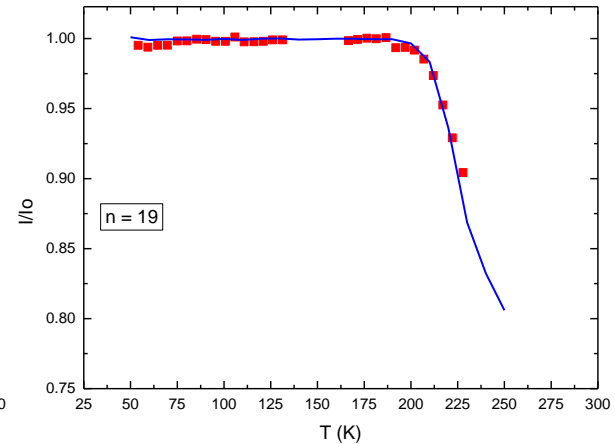
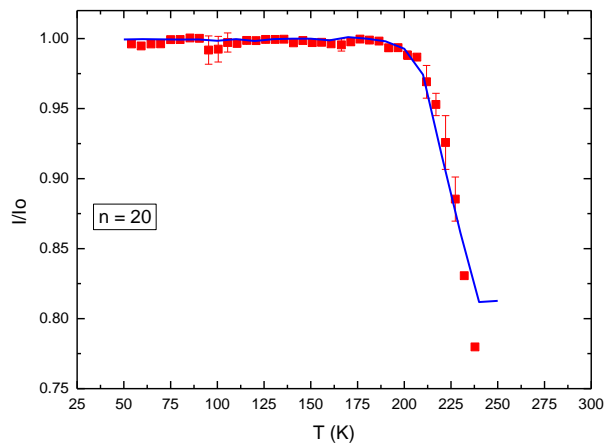
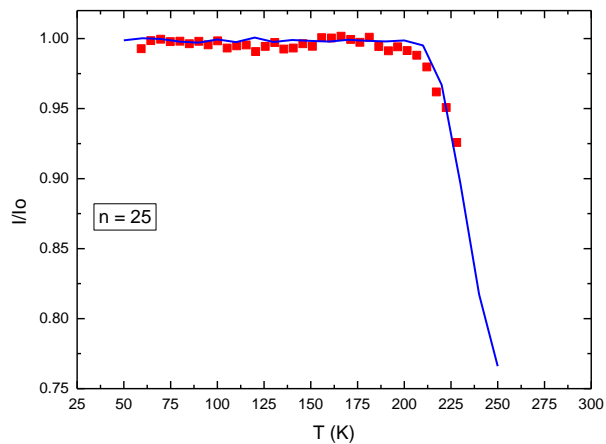
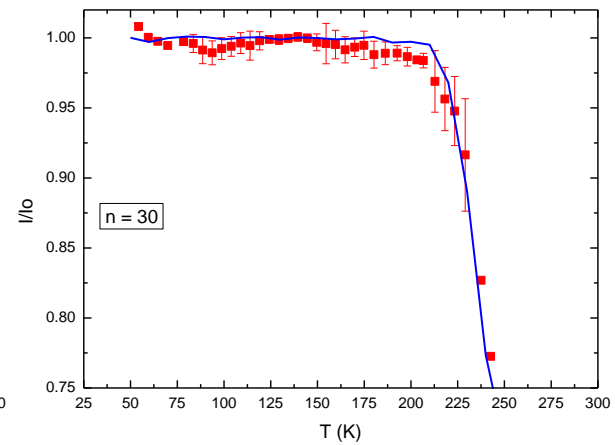
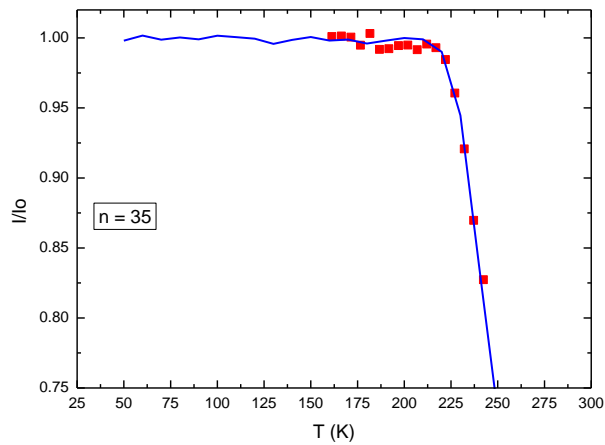
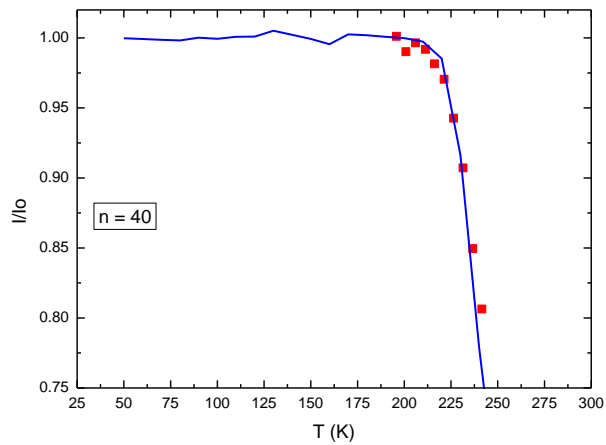


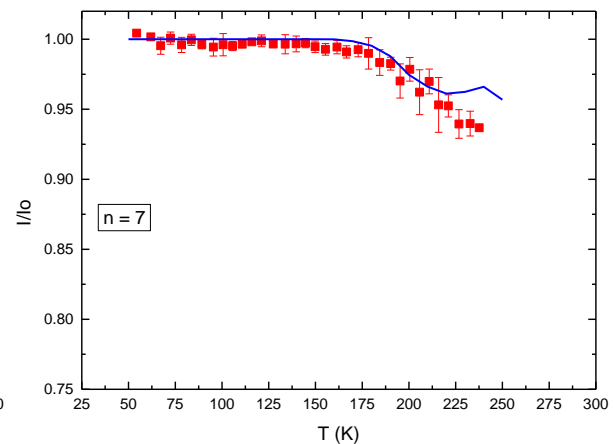
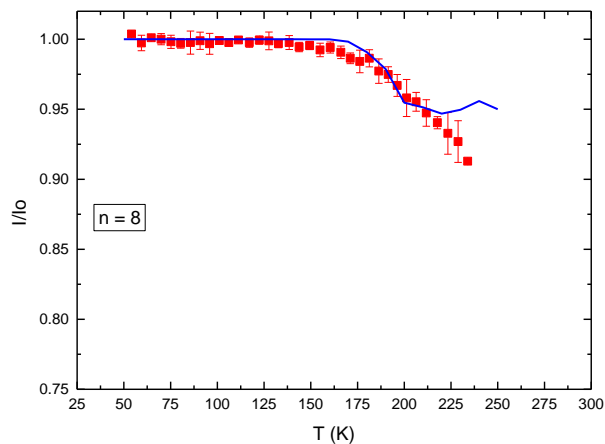
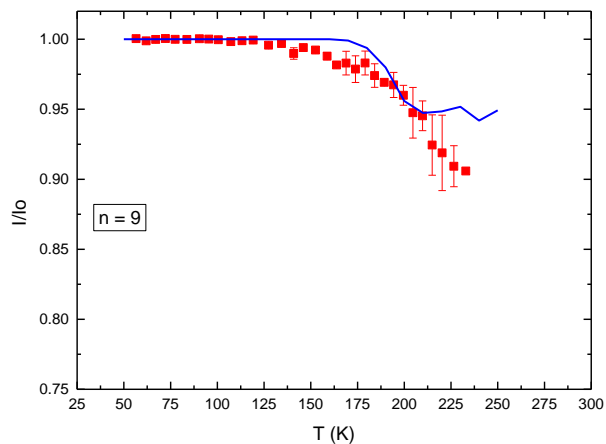
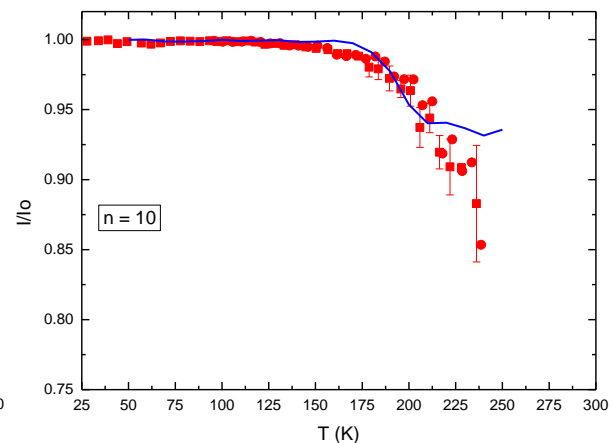
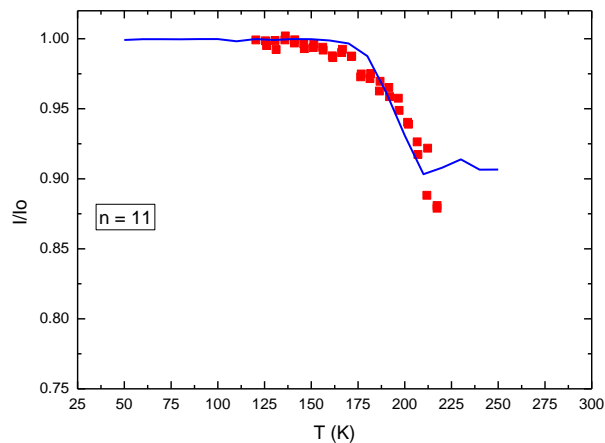
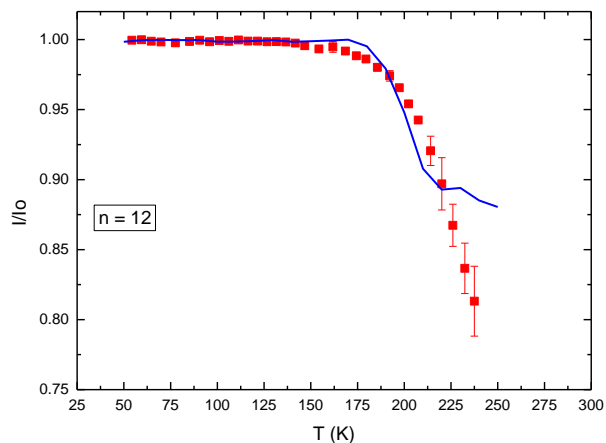
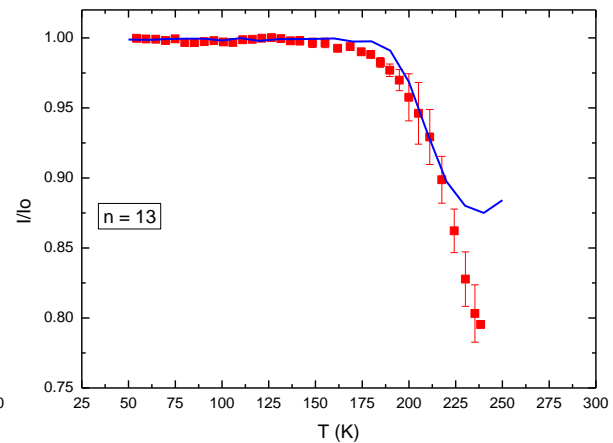
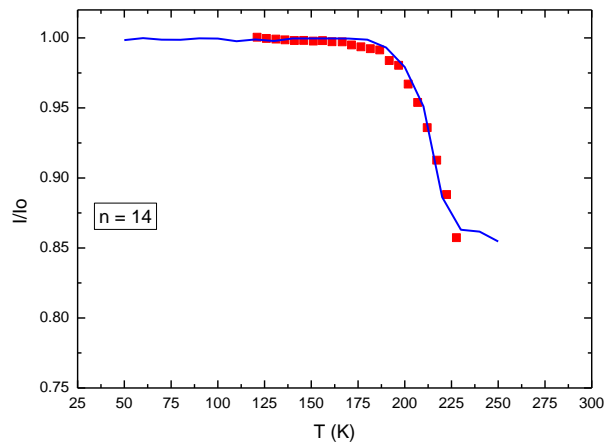
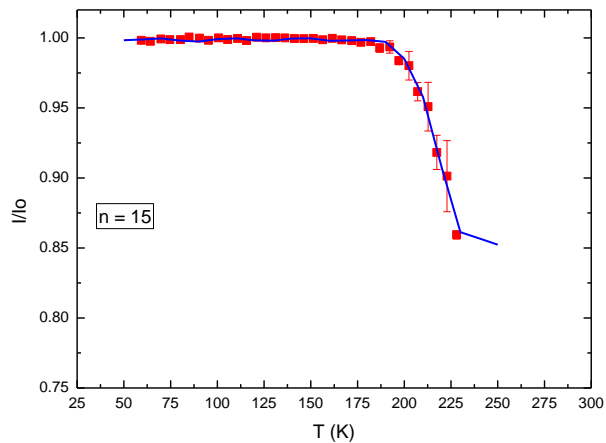
Generate initial population of size n (cascaded evaporations \Rightarrow use of dissociation energies of sizes $n+1, n+2, \dots$) (t_1)

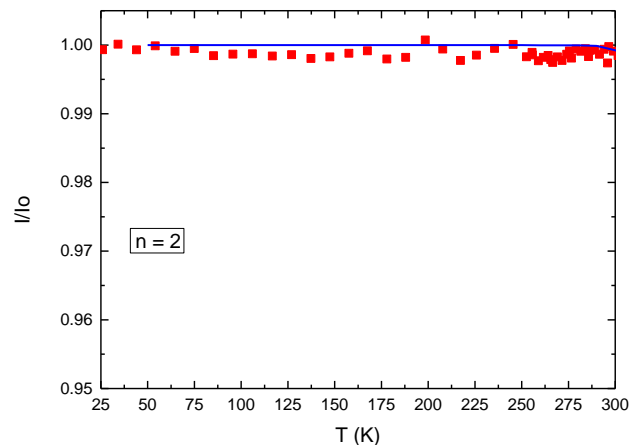
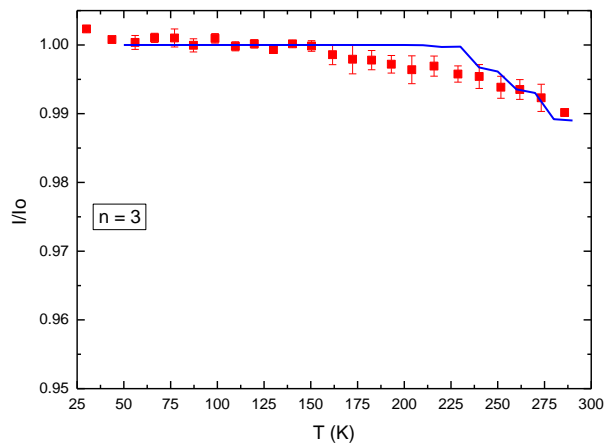
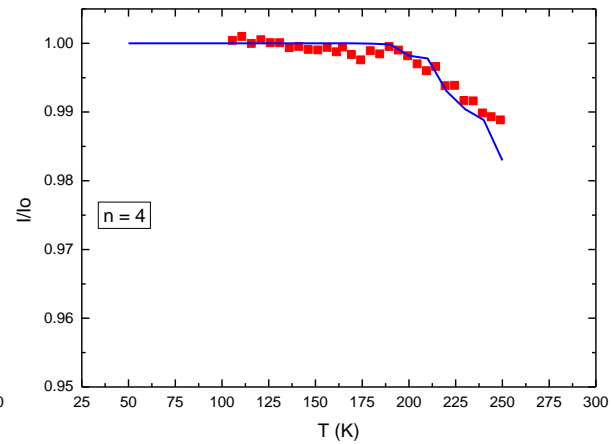
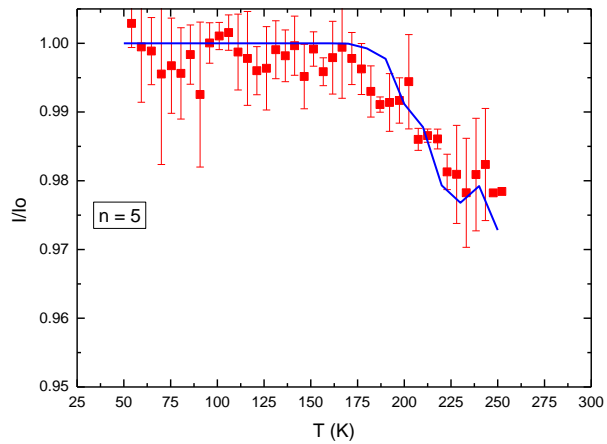
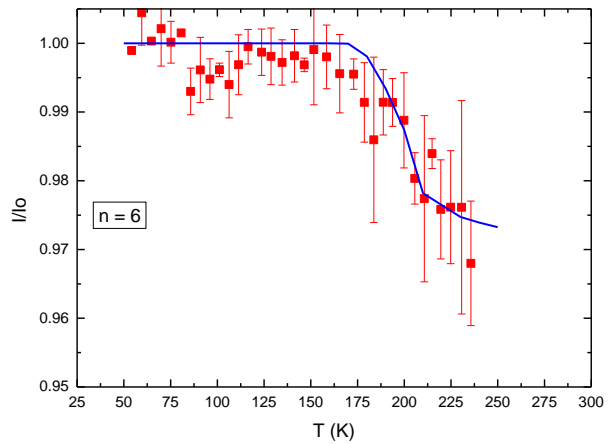
Calculate clusters trajectories with evaporation probabilities (t_2, t_3, t_4)

Generate TOF mass spectra *vs* initial temperature

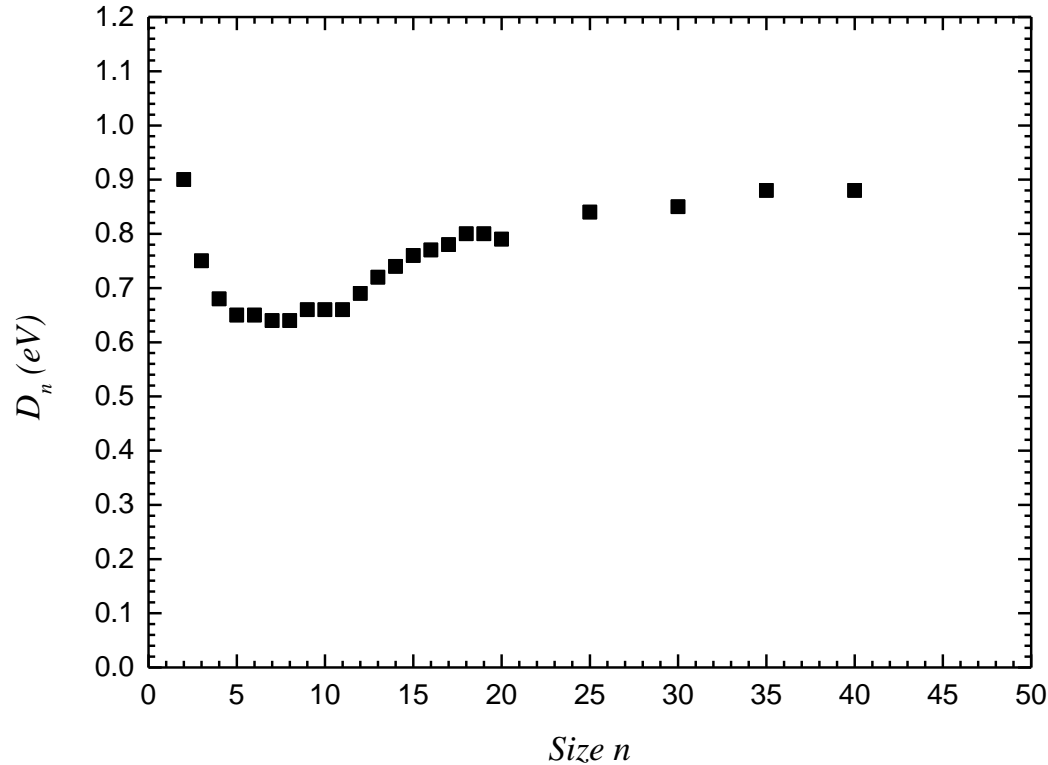
Compare with experiment, adjust dissociation energy





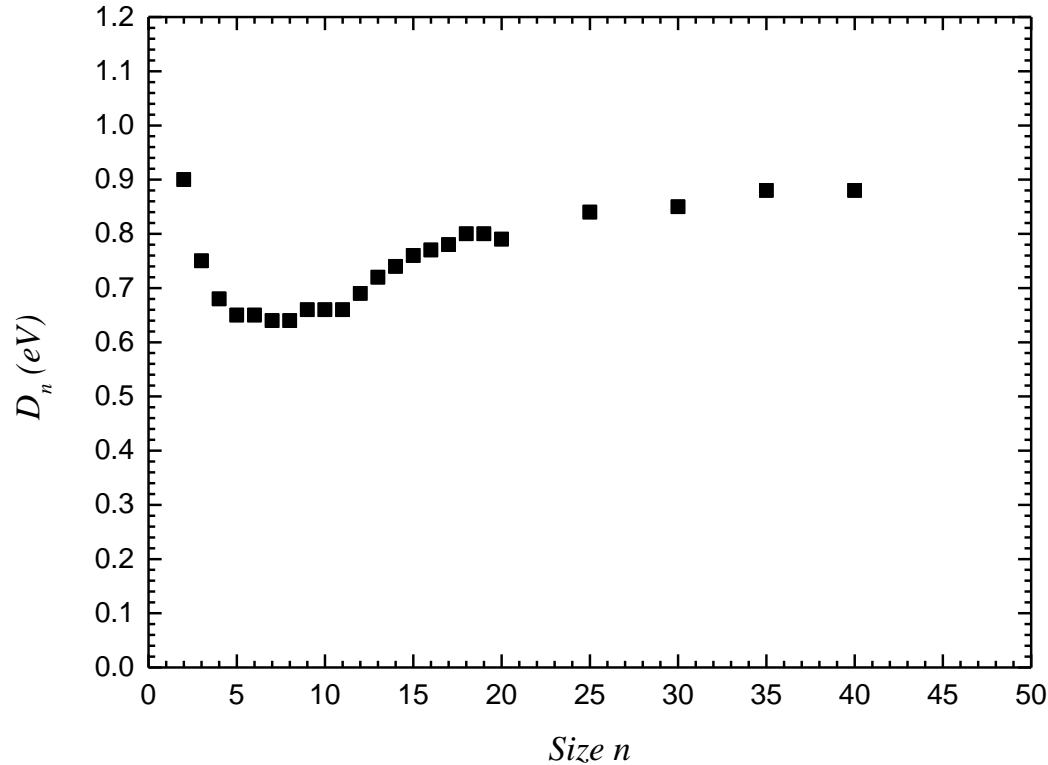


Dissociation energies



Black squares: values deduced from the reproduction of the experimental curves using PST

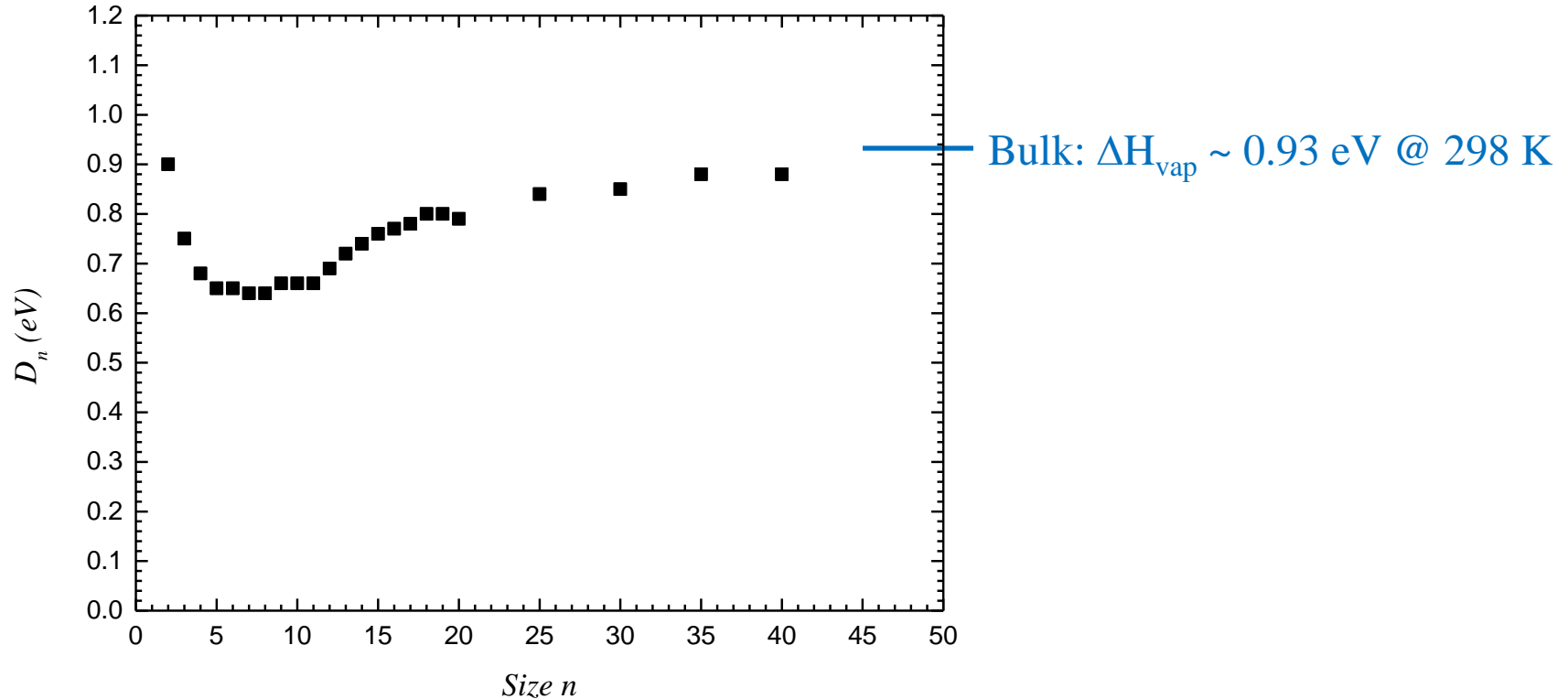
Dissociation energies



Black squares: values deduced from the reproduction of the experimental curves using PST

Only lower limit for $n=2$

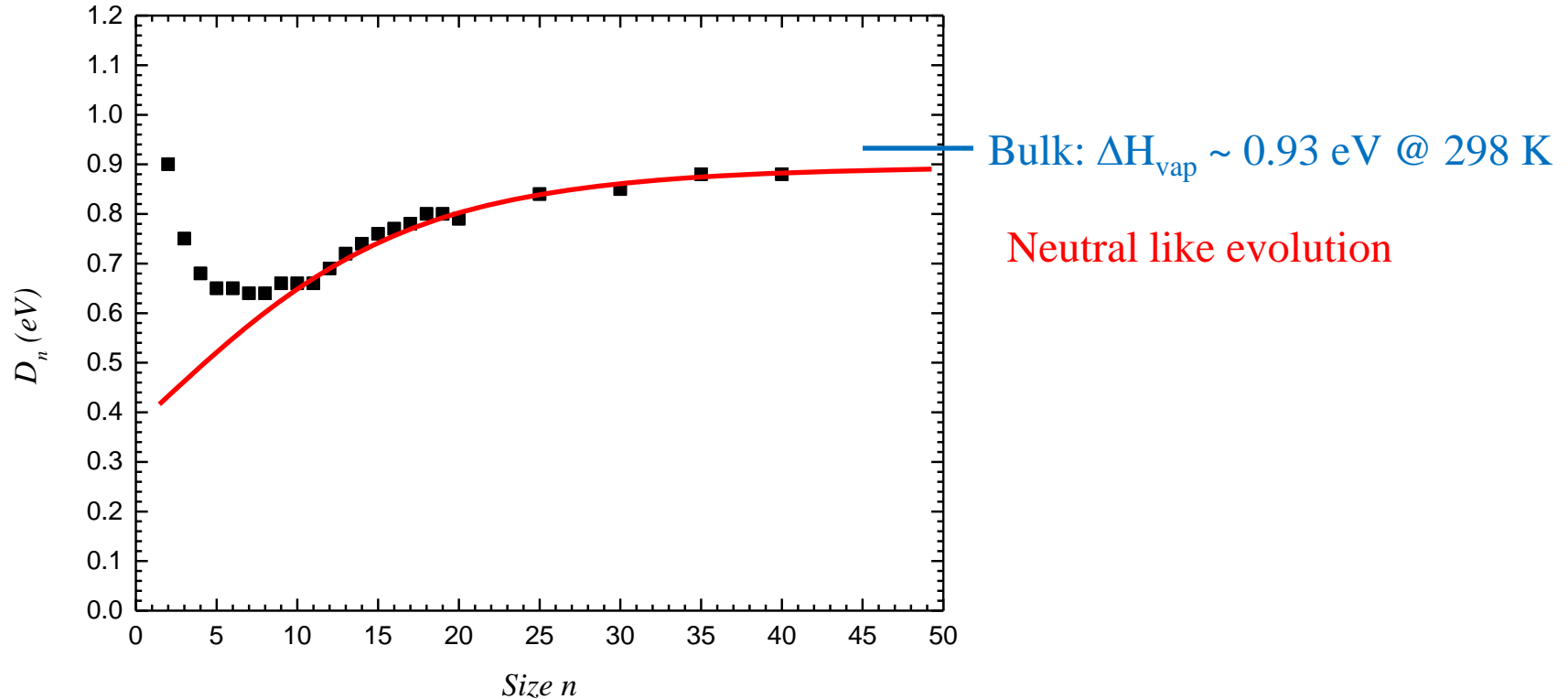
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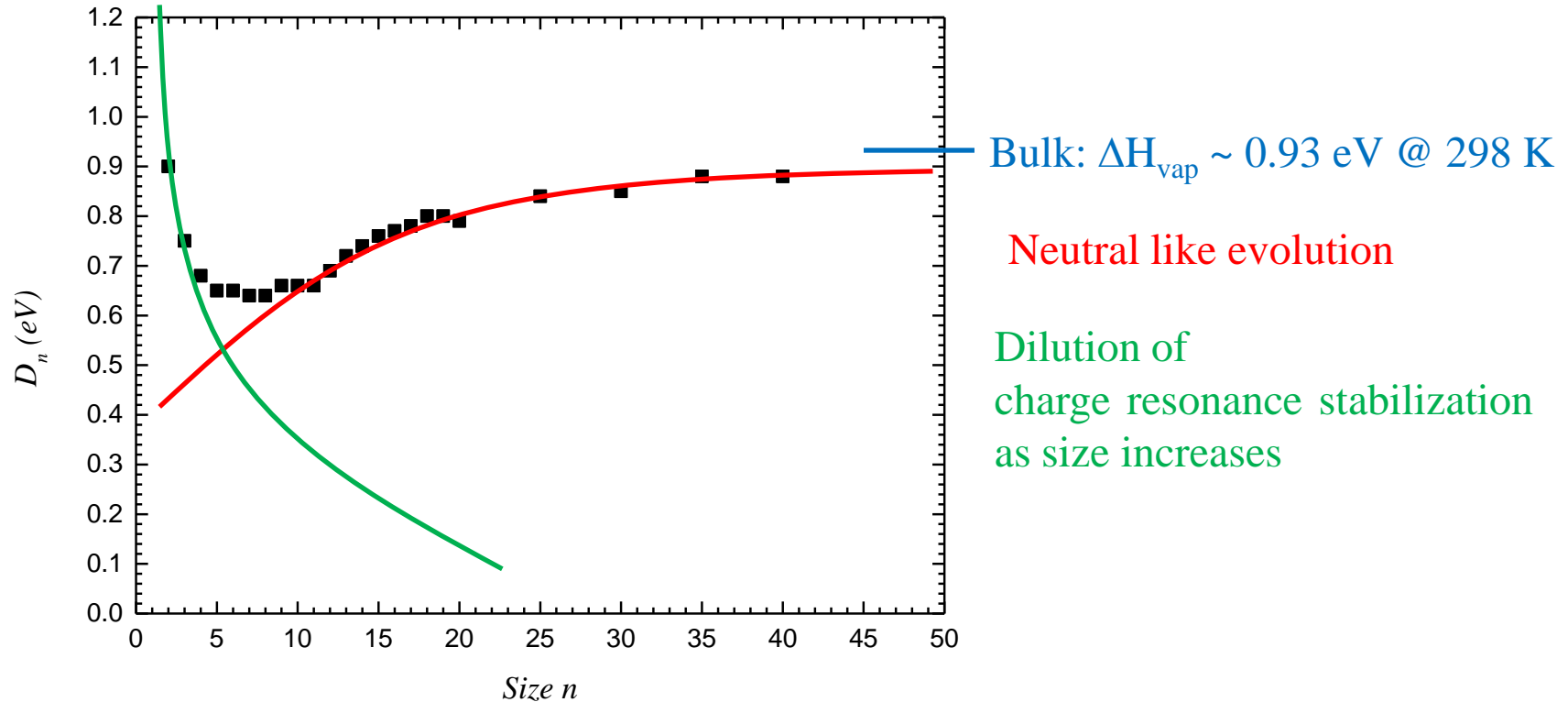
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Dissociation energies



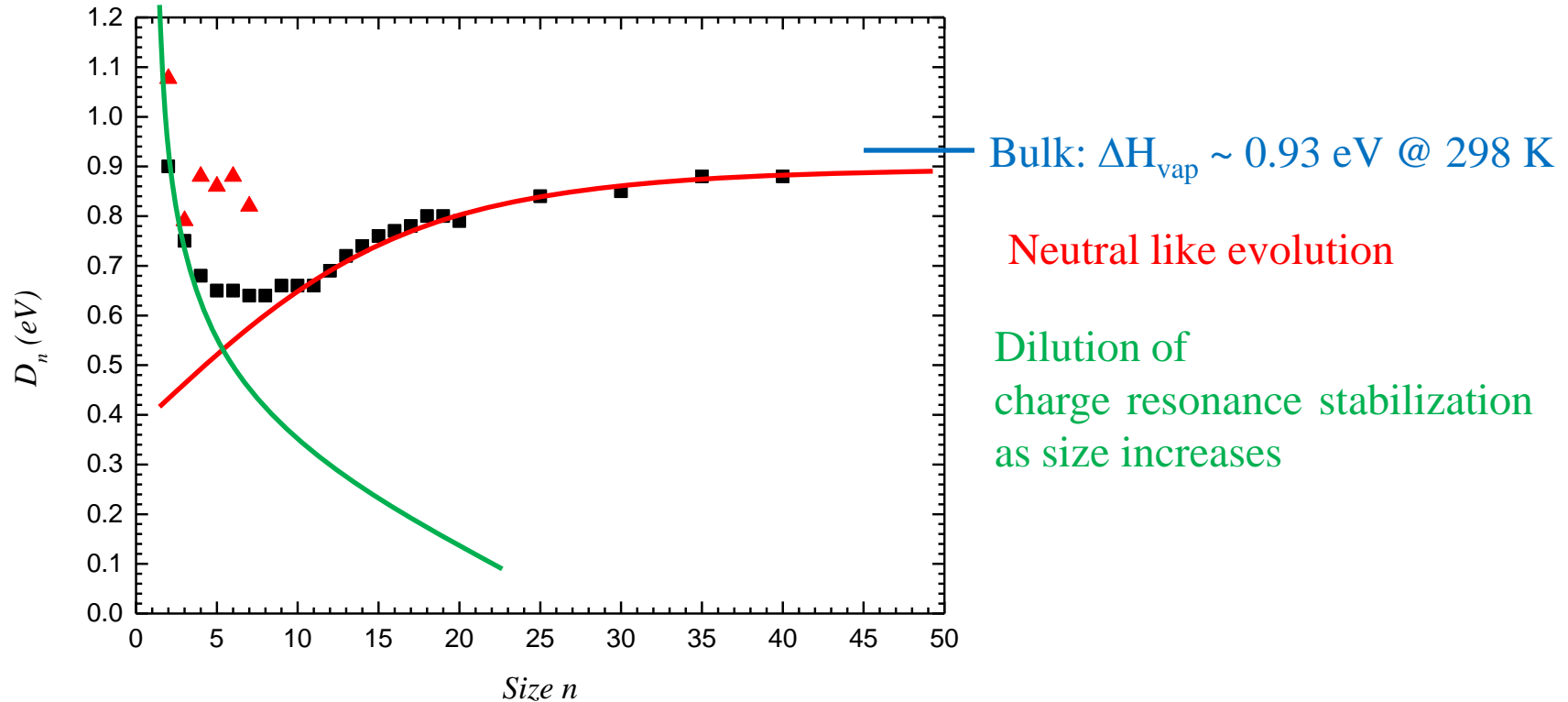
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Dissociation energies



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Dissociation energies



Black squares: values deduced from the reproduction of the experimental curves using PST

Only lower limit for $n=2$

Red triangles: DFTB calculation (cf Rapacioli *et al.*)

Conclusion

- We have observed the thermal evaporation of pyrene clusters, $n=2-40$
- Experimental results successfully reproduced using PST
- Dissociation energies deduced, in good agreement with theory and bulk value

Conclusion

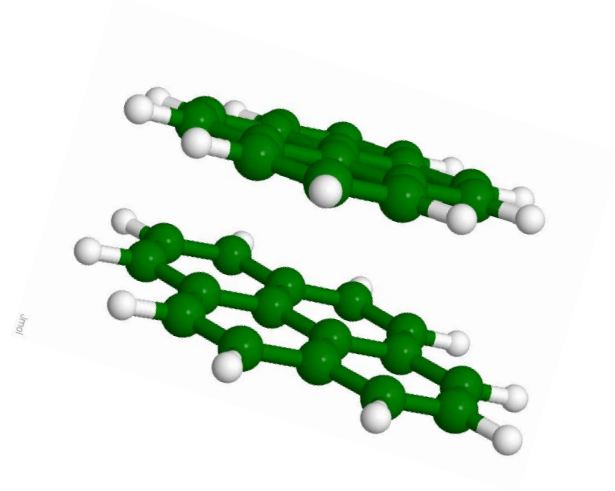
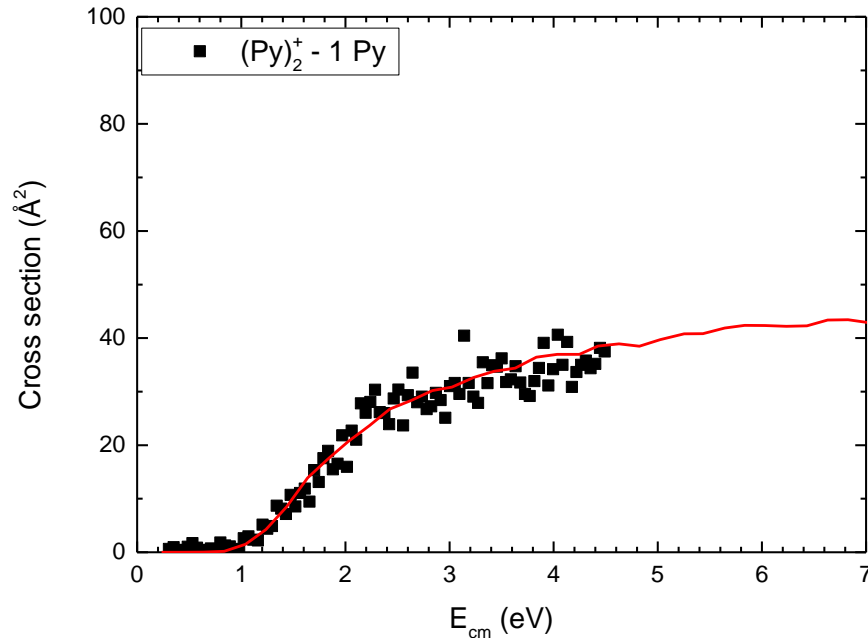
- Apart from the dissociation energy, no free parameters for the PST.

Too good to be true?

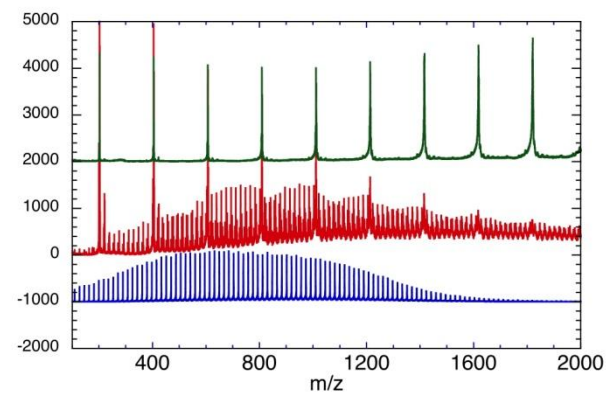
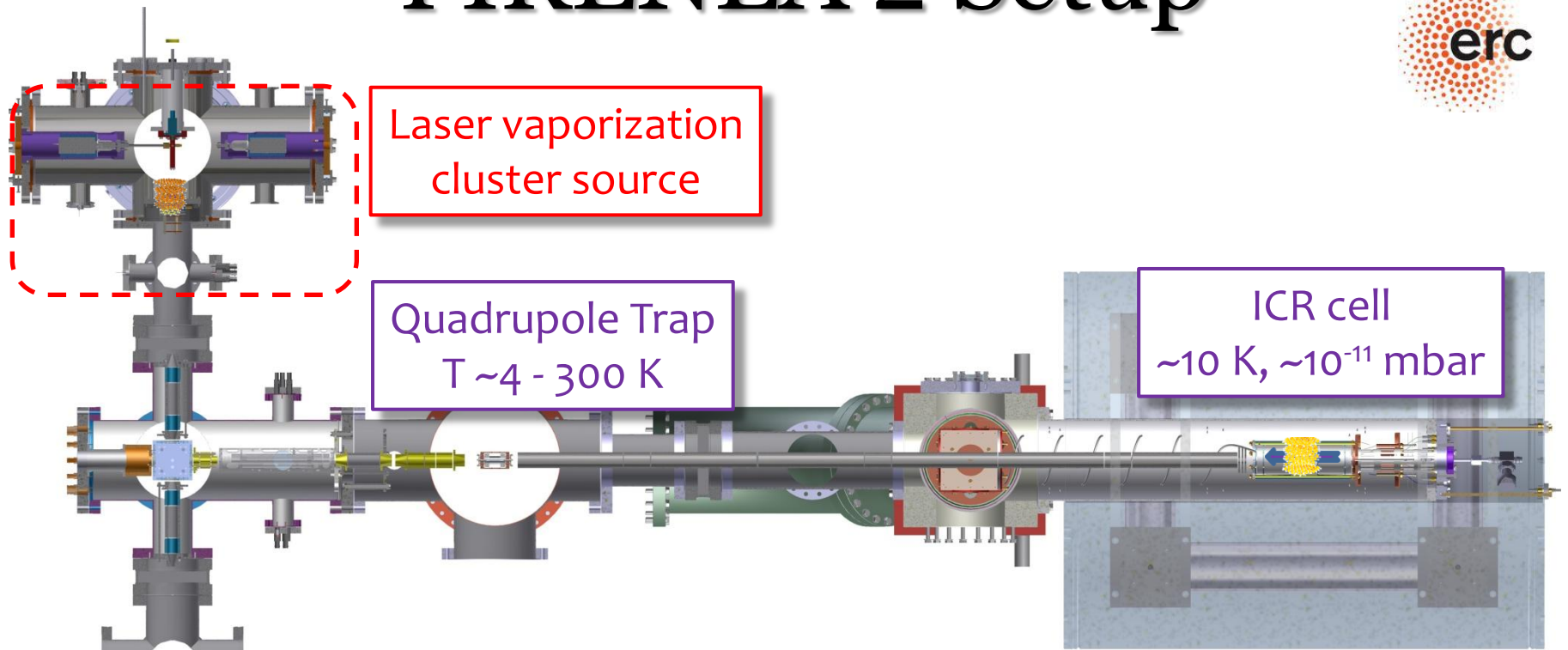
- At least, effective evaporation rates are obtained

Conclusion

- Preliminary data on CID cross-section measurements: dissociation energies compare well with evaporation data.

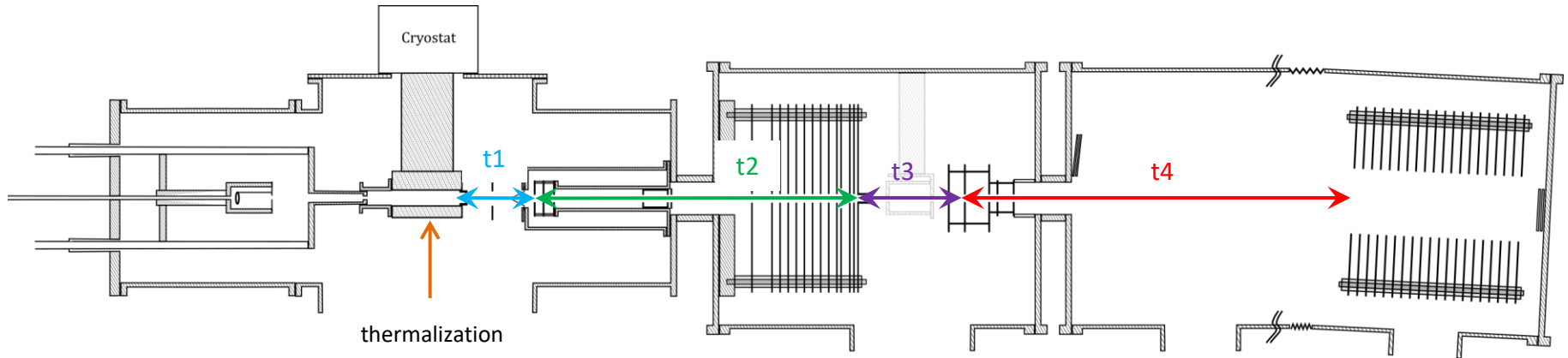


PIRENEA 2 Setup



$(\text{PAH})_n^+$
Mixed PAH-H₂O
aggregates
 $(\text{H}_2\text{O})_m\text{H}^+$

Experimental results



For a size n @ 22 eV, we have:

$$t_1 = D/v_0 = 20 \mu\text{s} @ 200 \text{ K}$$

$$t_2 = 24 \times \sqrt{n} \mu\text{s}$$

$$t_3 = 25 \times \sqrt{n} \mu\text{s}$$

$$t_4 = 25 \times \sqrt{n} \mu\text{s}$$

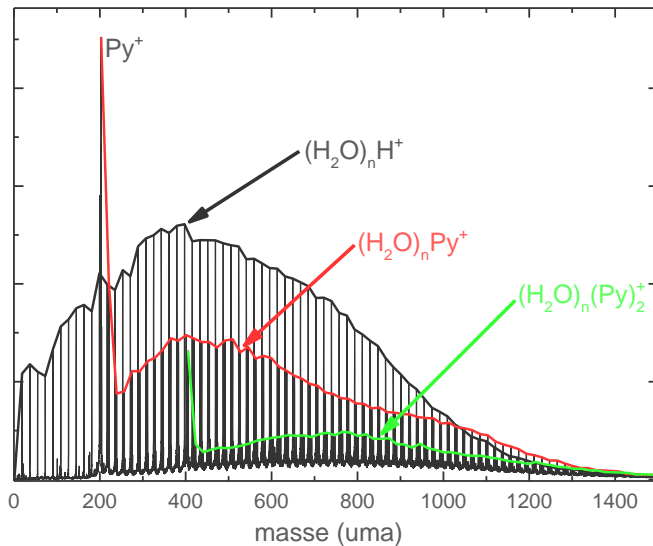
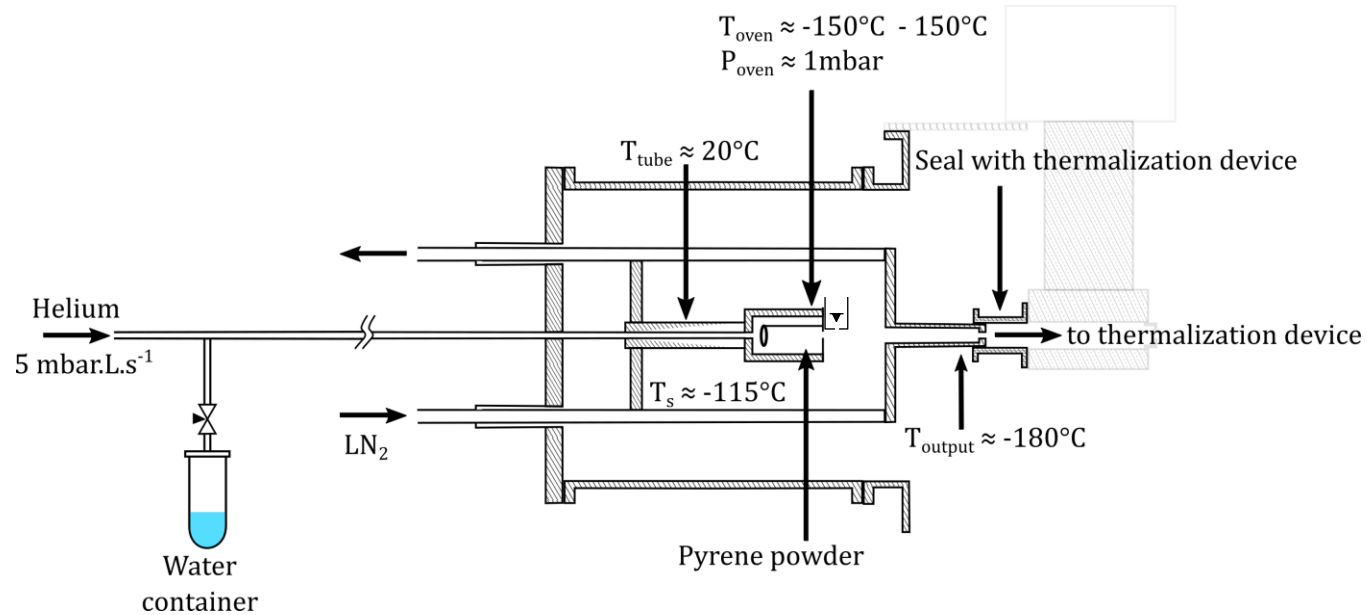
D = distance between the exit of the thermalizer and the first Wiley-McLaren = 8 cm

v_0 = thermal velocity of helium

Similar time spent in the different parts of the setup

⇒ Similar evaporation probability at each experimental stage

Cluster production



Small pure Pyrene clusters.

Large pure pyrene clusters

Mixed water-pyrene clusters

PST evaporation rates

The version of PST used here contains the constraint that total angular momentum is conserved, and that in the reverse process the associating product must overcome the maximum of the centrifugal barrier which defines the position of the transition state (loose transition state).

The interaction between the products is represented by an effective central potential of the form:

$$V_{eff}(r) = -\frac{C_4}{r^4} + \frac{L^2 \hbar^2}{2\mu r}$$

PST evaporation rates

Ingredients of the Phase Space Theory:

- initial internal energy of the parent E_i
- density of states of the parent $N(E)$
- total number of states of the fragments $G(E, J)$
- conservation of angular momentum

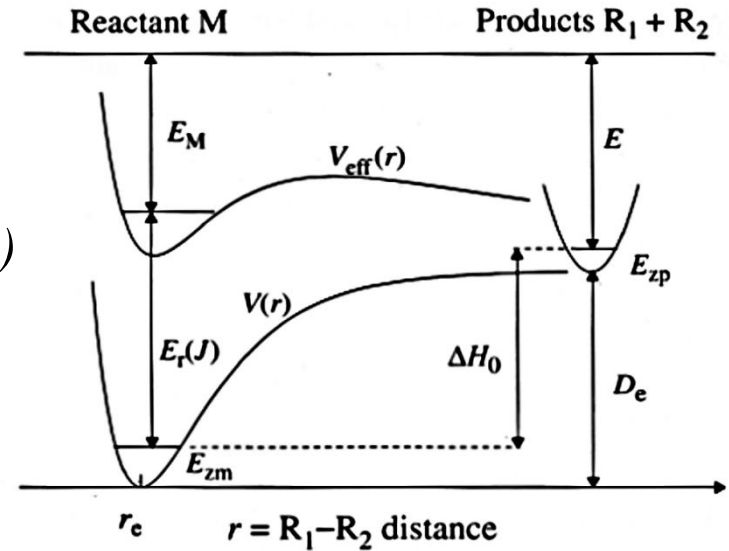
$$W(E_i, J) = \frac{G(E_f, J)}{h(2J + 1)N(E_i)}$$

$$E_f = E_i + E_{rot} - D$$

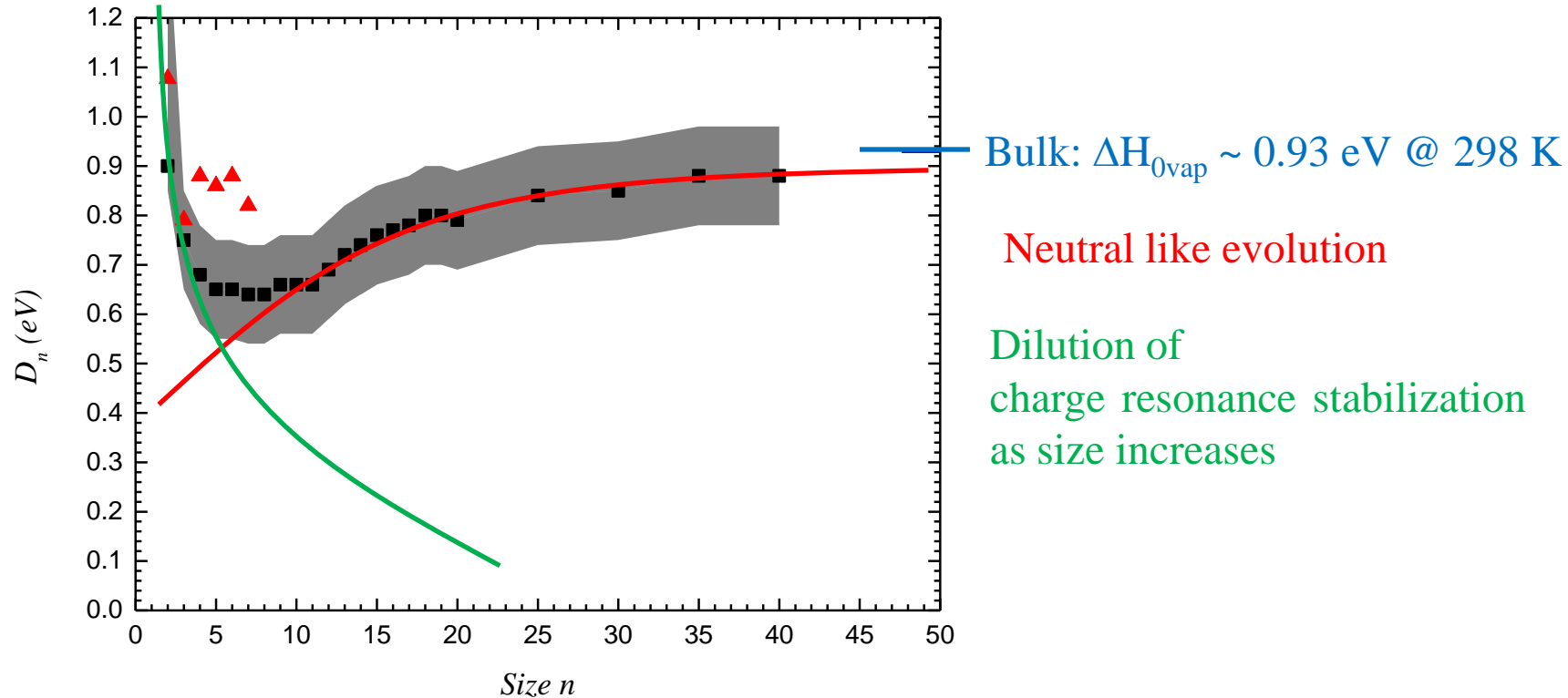
$$E_{rot} = B_0 J(J + 1)$$

From the PST we also get the probability of energy partition among the fragments:

$$P(E_v, E_{rel}, E_r) dE_v dE_{rel} dE_r = \frac{N_1(E_v) N_2(E_f - E_v - E_{rel} - E_r) N_x(E_r, E_{rel}, J)}{G(E_f, J)} dE_v dE_{rel} dE_r$$



Dissociation energies

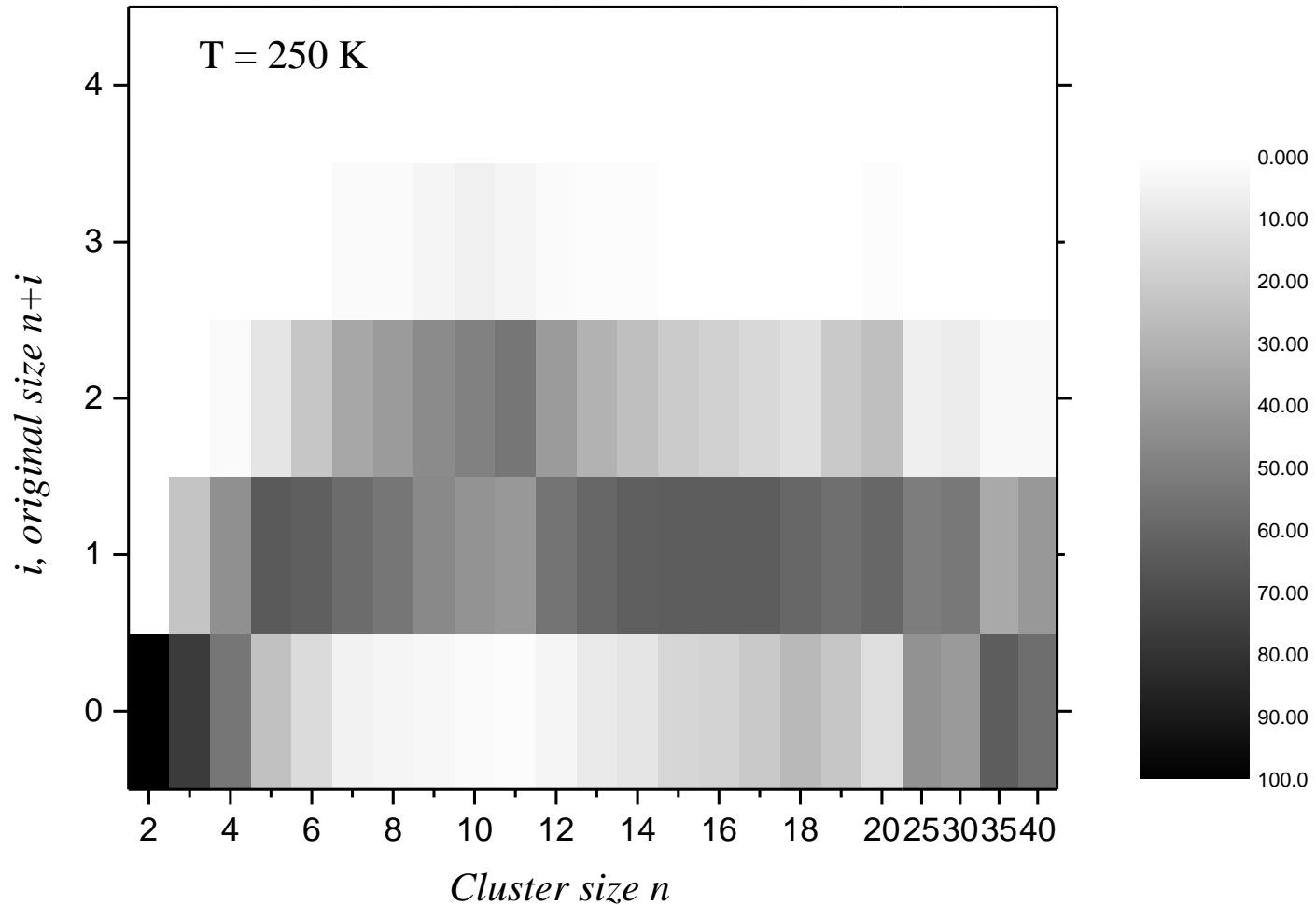


Black squares: values deduced from the reproduction of the experimental curves using PST

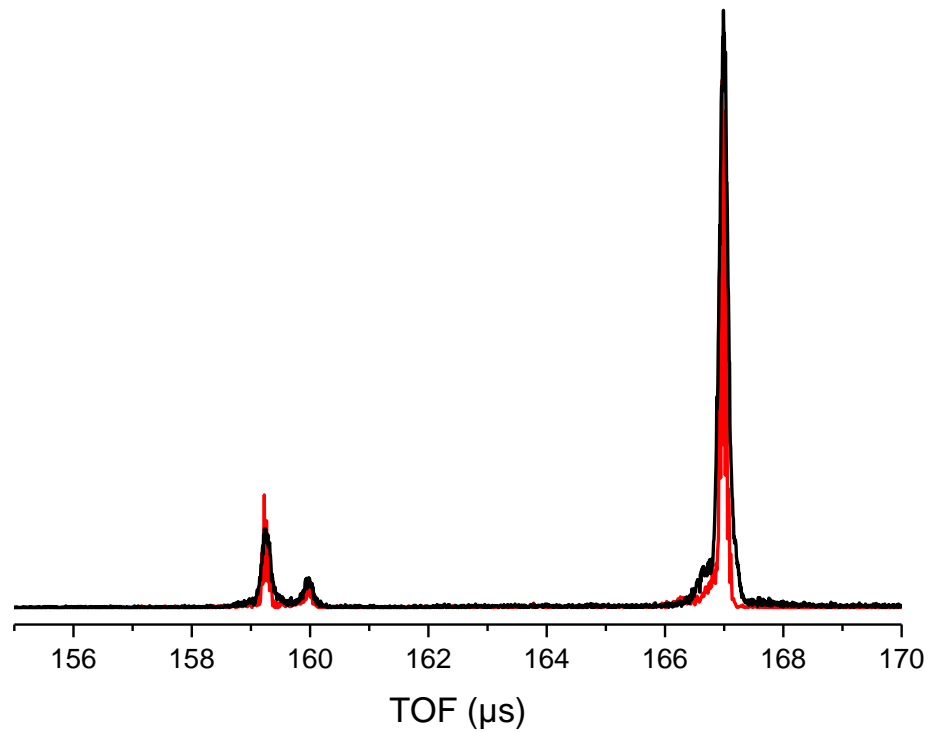
Red triangles: DFTB calculation (cf Rapacioli *et al.*)

Bulk: $\Delta H_{0\text{vap}} \sim 0.93$ eV @ 298 K

Initial size contributions



Simulated TOF mass spectra



Monte Carlo integration

To compare with the experiment we need to integrate over distributions:

$$\sum_J \int P(E,T)P(J,T)dE$$

$P(E,T)$ = probability to have an internal energy E in the parent cluster.

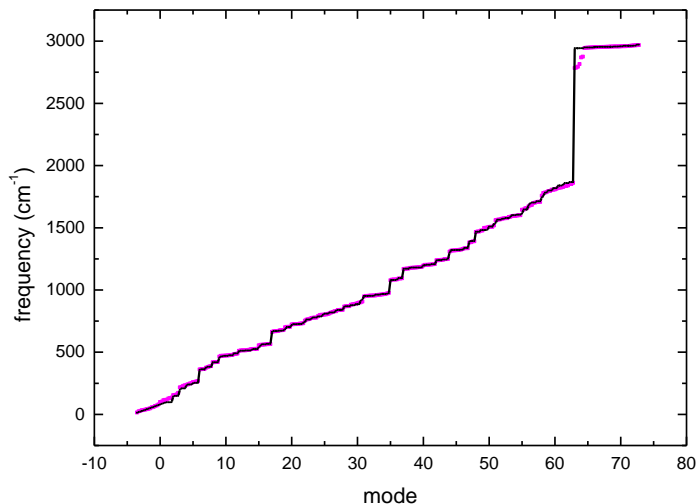
Number occupation n_i of each vibration randomly picked up so that on average we have:

$$E = \sum_i \frac{\hbar\omega_i}{e^{\frac{\hbar\omega_i}{k_B T}} - 1}$$

$P(J,T)$ = probability to have the total angular momentum J in the parent cluster.

$$P(J) = \frac{4}{\sqrt{\pi}} \left(\frac{B}{k_B T} \right) J^2 e^{-\frac{BJ^2}{k_B T}} \quad \text{High temperature or low B limit used}$$

Vibrational harmonic frequencies



harmonic frequencies from DFTB for the tetramer

Vibrational harmonic frequencies calculated for neutral pyrene and cationic clusters of sizes 1, 2, 3 and 4.

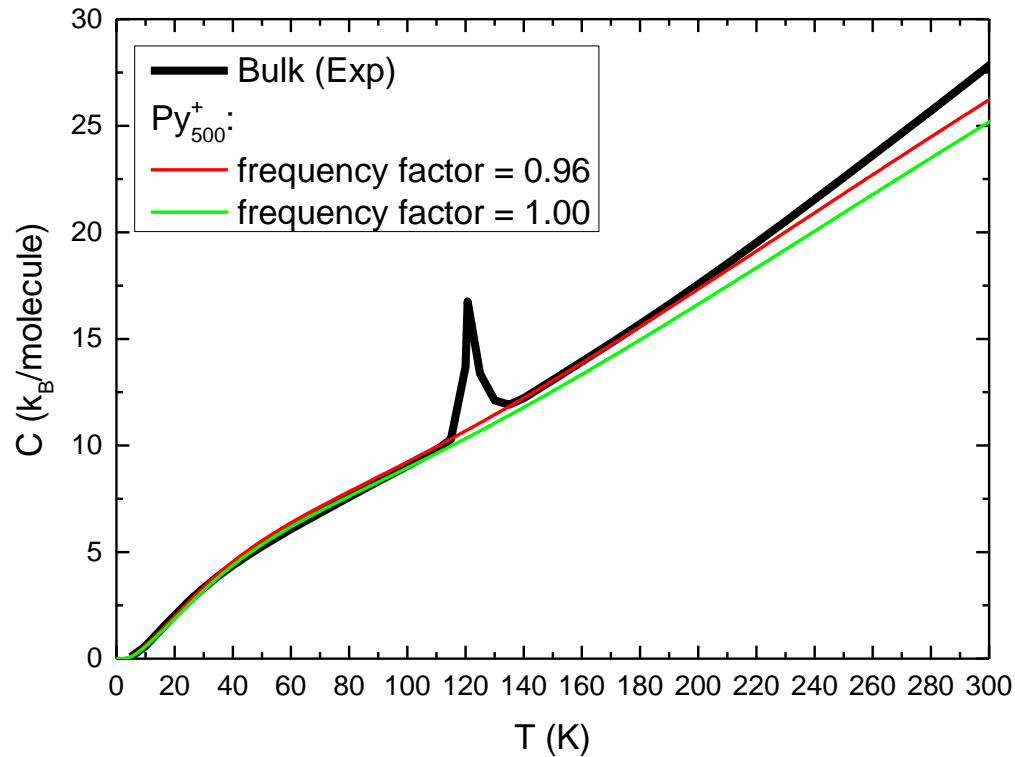
To extrapolate for larger sizes n :

Intramolecular frequencies = cation + $(n-1)$ neutrals

Intermolecular frequencies = linear interpolation between 13 cm⁻¹ and the lowest cationic frequency.

Vibrational harmonic frequencies

Using these, the heat capacity of bulk pyrene can be reproduced:

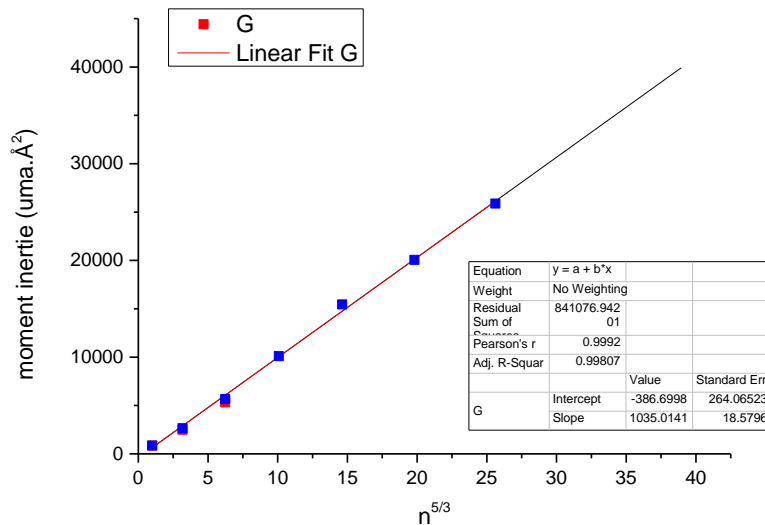


Moments of inertia and rotational constants

For the PST calculation the rotational constant B is needed. $B = \frac{\hbar^2}{2I}$

The moment of inertia are taken as the geometrical average of the 3 main moments:

$$I = \sqrt[3]{I_1 I_2 I_3}$$

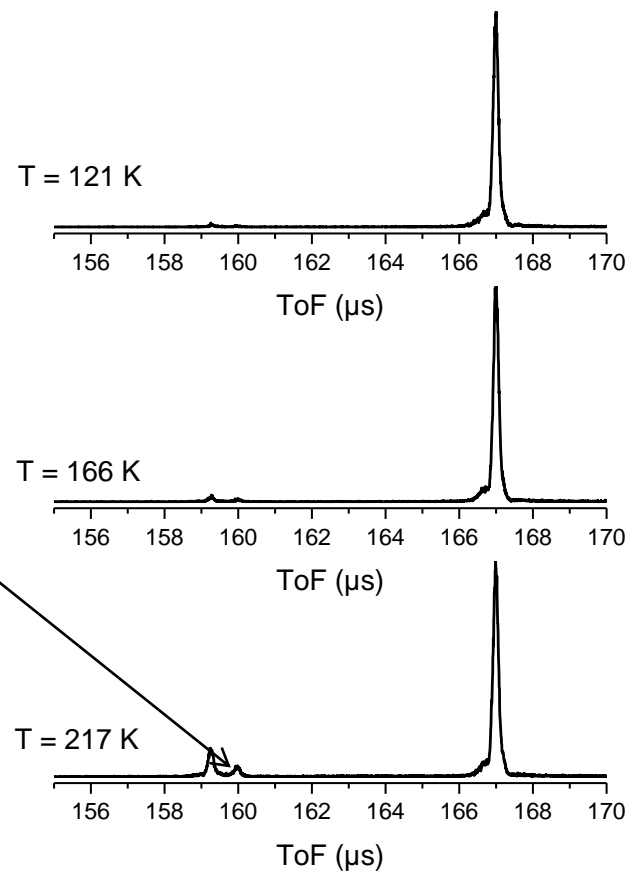


Squares: DFTB calculation

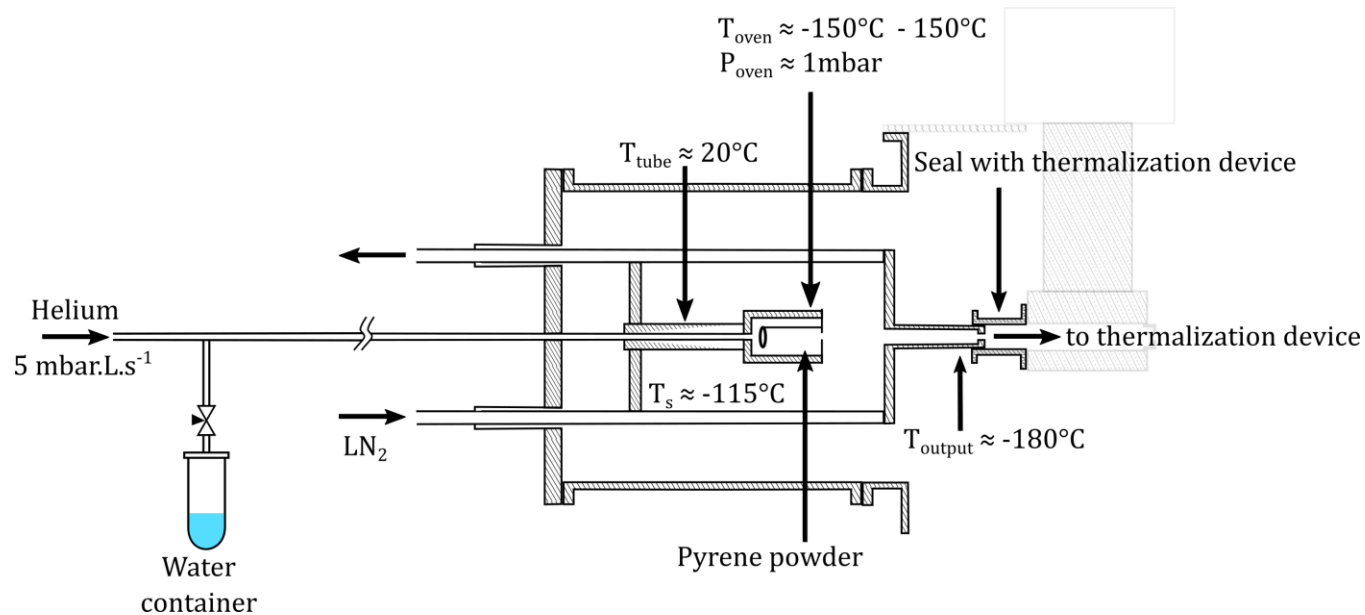
Line : $I(Py_n) \propto I(Py_1)n^{5/3}$

If clusters evaporate after the second Wiley McLaren acceleration and before the reflectron (t_4), they end up in the “secondary” fragment peak.

If they evaporate after the reflectron they can not be distinguished from the parent peak.



Clusters production

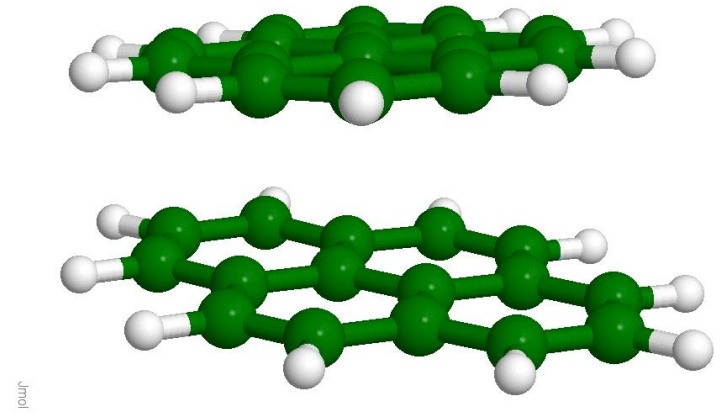
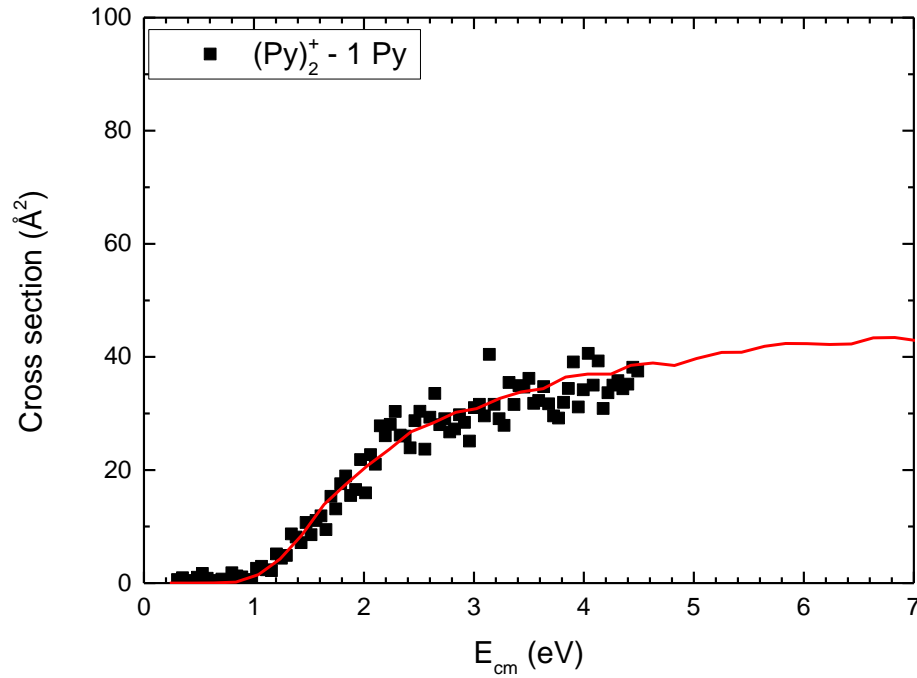


Clusters are produced in a gas aggregation source.

Ionization by either :

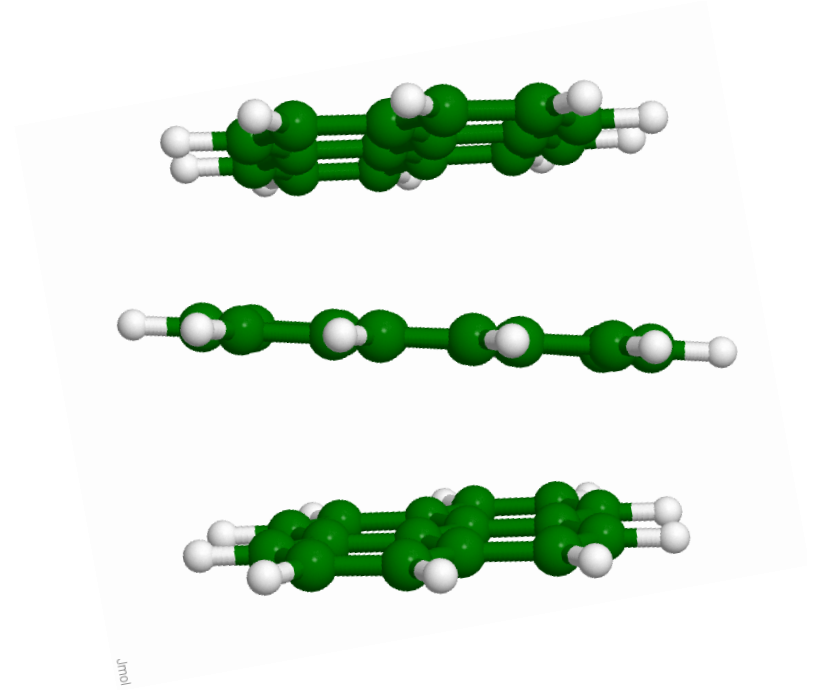
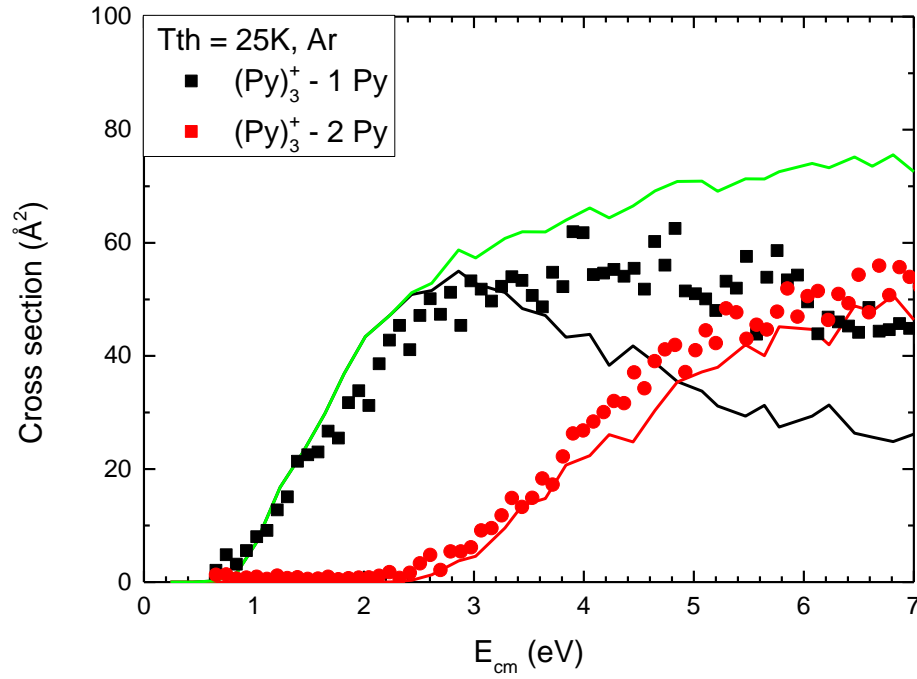
- a discharge electrode, $I \approx 100\mu\text{A}$, $V \approx 1 \text{ kV}$
- an electron gun, $I = 2.5 \text{ A}$, $V = -100 \text{ V}$

Experimental results

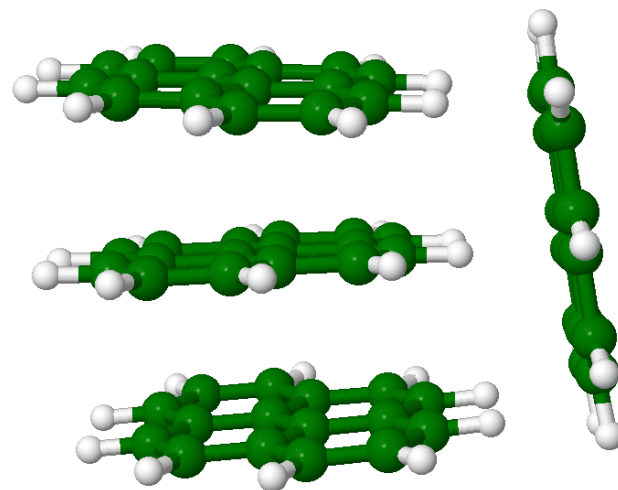
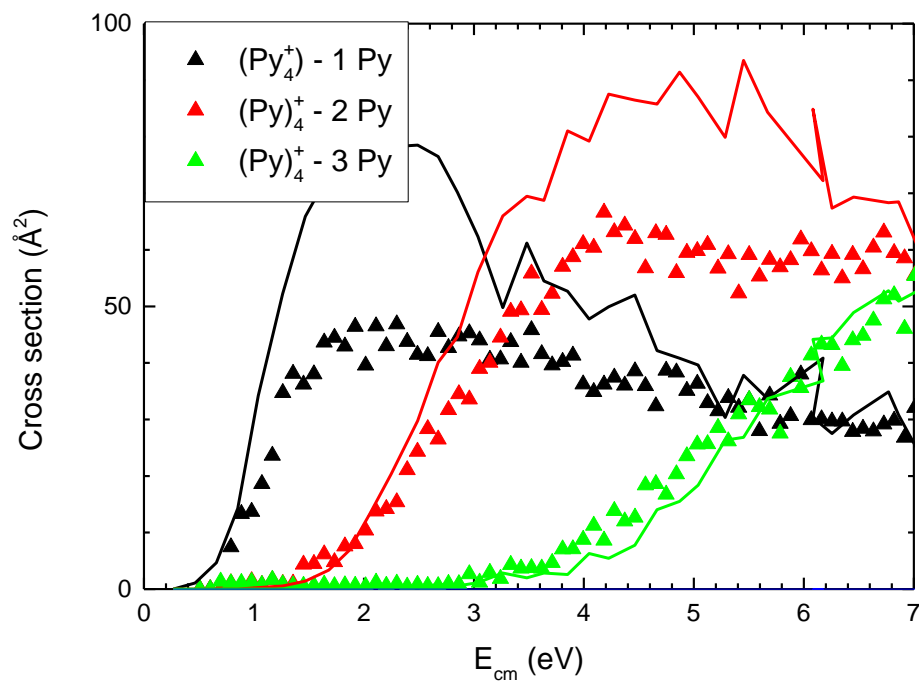


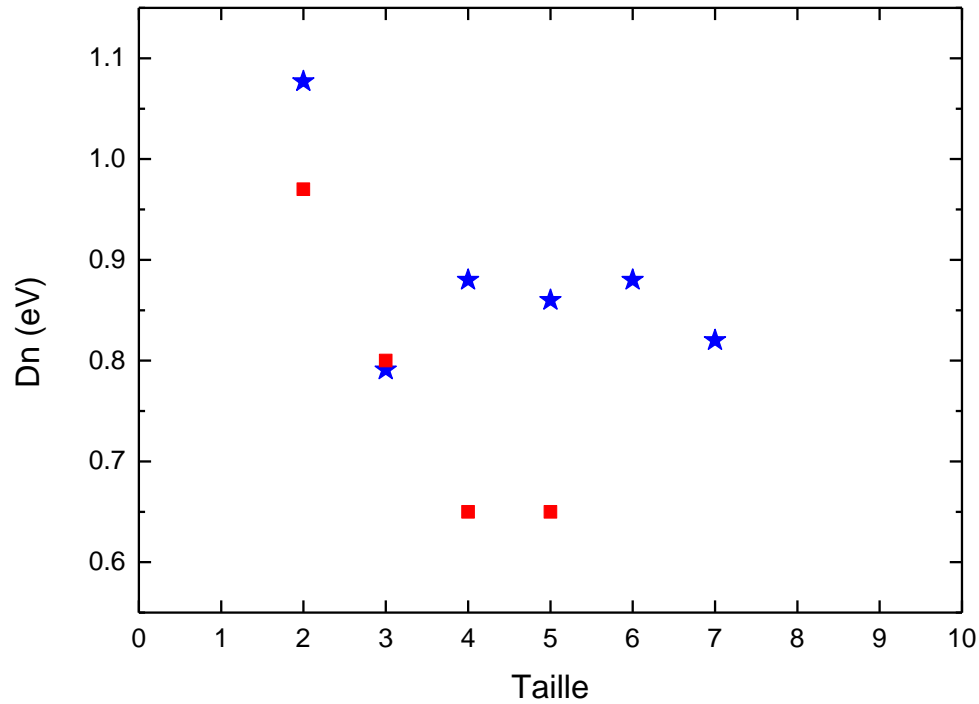
pyr

Experimental results



Experimental results





Blue stars: theory

Red squares: values deduced from the reproduction of the experimental curves using PST (CID).